

## LOSS OF EXPONENTIAL STABILITY FOR A THERMOELASTIC SYSTEM WITH MEMORY

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ABSTRACT. In this article we study a thermoelastic system considering the linearized model proposed by Gurtin and Pipkin [8] instead of the Fourier's law for the heat flux. We use theory of semigroups [9, 11] combining Pruss' Theorem [10] and the idea developed in [5] to show that the system is not exponentially stable.

### 1. INTRODUCTION

We study a partial differential equation that models an elastic string:

$$u_{tt} - u_{xx} + \theta_{xx} = 0 \quad \text{in } (0, L) \times (0, \infty), \quad (1.1)$$

$$\theta_t - g * \theta_{xx} + c g * \theta - u_{xxt} = 0 \quad \text{in } (0, L) \times (0, \infty), \quad (1.2)$$

with initial data

$$u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), \quad \theta(x, 0) = \theta_0(x).$$

The function  $u = u(x, t)$  is the small transversal vibration of the elastic string of reference configuration of length  $L$ , and  $\theta = \theta(x, t)$  is the temperature difference from the material and natural ambient. To fix ideas we assume that the string is held fixed at both ends,  $x = 0$  and  $x = L$ . We impose the boundary conditions

$$u(0, t) = u(L, t) = 0,$$

$$\theta(0, t) = \theta(L, t) = 0.$$

In this model,  $c$  is a positive constant, and  $g : \mathbb{R}^+ \rightarrow \mathbb{R}^+$  is the relaxation function. We assume that  $g$  is differentiable and satisfies  $g(0) > 0$ ,  $g'(t) < 0$  and

$$1 - \int_0^\infty g(s) ds = \ell > 0.$$

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We introduce the convolution product

$$(g * u)(t) := \int_0^t g(t - \tau)u(\cdot, \tau)d\tau.$$

Now we observe that when  $c = 0$  the thermoelastic system has exponential decay, as can be seen in [4], when we replace  $g * u$  by  $\theta$  in (1.2) we also have exponential decay, see [3]. The similar situation is valid for thermoelastic plate, see [5] and [7].

The article is organized as follows, in the Section 2 we introduce the notation and the functional spaces, in the Section 3 we obtain the semigroup of solutions and finally, in the Section 4 we prove the loss of exponential stability for the thermoelastic system with memory.

## 2. FUNCTIONAL SETTING AND NOTATION

We use the standard Lebesgue spaces and Sobolev spaces with their usual properties as in [1]. Consider the positive operators  $A$  and  $B$  on  $L^2(0, L)$  defined by  $A = -(\cdot)_{xx}$  and  $B = cI - (\cdot)_{xx}$ , with domains  $D(A) = D(B) = (H^2 \cap H_0^1)(0, L)$ . Now, for  $r \in \mathbb{R}$ , we introduce the scale of Hilbert spaces  $H_r = D(A^{r/2})$  with the usual inner products  $\langle v_1, v_2 \rangle_{H_r} = \langle A^{r/2}v_1, A^{r/2}v_2 \rangle$  and we have  $H_{r_1} \hookrightarrow H_{r_2}$  are compact whenever  $r_1 > r_2$ . Concerning the memory kernel  $g$ , we make the substitution  $\mu(s) = -g(s)$  and we require

$$\mu \in C^1(\mathbb{R}^+) \cap L^1(\mathbb{R}^+), \quad \mu(s) > 0, \quad \mu'(s) \leq 0, \quad g(0) = \int_0^\infty \mu(s)ds > 0. \quad (2.1)$$

Calling  $\sigma_\infty = \sup\{s : \mu(s) > 0\}$ , we infer that, dual to (2.1), for each  $\sigma > 0$ , there exists a set  $\mathcal{O}_\sigma \subset (\sigma, \sigma_\infty)$  of positive Lebesgue measure such that  $\mu'(s) < 0$ , in  $\mathcal{O}_\sigma$ . Now for  $r \in \mathbb{R}$  consider the weighted Hilbert spaces:

$$\mathcal{M}_r = L_\mu^2(\mathbb{R}^+; H_r)$$

with the inner product

$$\langle \nu, \eta \rangle_{\mathcal{M}_r} = \int_0^\infty \mu(s) \langle B^{r/2}\nu(s), B^{r/2}\eta(s) \rangle ds \quad (2.2)$$

and we introduce as in [6] the linear operator  $T$  on  $\mathcal{M}_1$  defined by  $T\eta = -\eta_s$  with domain

$$D(T) = \{\eta \in \mathcal{M}_1 : \eta_s \in \mathcal{M}_1, \eta(0) = 0\},$$

where  $\eta_s$  is the distributional derivative of  $\eta$  with respect to the internal variable  $s$ , and then the operator  $T$  is the infinitesimal generator of a  $C_0$ -semigroup of contractions. In particular, there holds

$$\langle T\eta, \eta \rangle_{\mathcal{M}_1} = \int_0^\infty \mu'(s) \|B^{1/2}\eta(s)\|^2 ds \leq 0, \quad \text{for all } \eta \in D(T). \quad (2.3)$$

Finally, we define with the usual inner products, the following Hilbert spaces:

$$\mathcal{H}_r = H_{r+2} \times H_r \times H_r \times M_{r+1}, \quad r \in \mathbb{R}.$$

## 3. THE SEMIGROUP OF SOLUTIONS

To describe properly the solutions of the system (1.1)-(1.2) by means of a  $C_0$ -semigroup of linear operators acting on the phase-space  $\mathcal{H}_0$ , we will follow the ideas of [1]. In this direction we introduce an additional variable, namely, the summed past history of  $\theta$  defined as

$$\eta^t(s) = \int_0^s \theta(t-y)dy, \quad \text{with } t, s \geq 0.$$

Observe that we have formally  $(\frac{d}{dt} + \frac{d}{ds})(\eta^t(s)) = \theta$  in  $(0, L)$  subject to the boundary and initial conditions  $\eta^t(0) = 0$  in  $(0, L)$ ,  $t \geq 0$ ,

$$\eta^0(s) = \int_0^s \theta(-y)dy, \quad s \geq 0.$$

For the rest of this article, we consider the vectors  $U(t) = (u(t), v(t), \theta(t), \eta^t)^T$  and  $U(0) = (u_0, v_0, \theta_0, \eta_0)^T \in \mathcal{H}_0$ . We obtain the linear evolution equation, in  $\mathcal{H}_0$ ,

$$U_t - LU = 0 \tag{3.1}$$

$$U(0) = U_0 \tag{3.2}$$

where the linear operator  $L$  is defined as

$$LU = \begin{pmatrix} v \\ u_{xx} - \theta_{xx} \\ u_{xx} - \int_0^\infty g(s)[c\theta(t-s) - \theta_{xx}(t-s)]ds \\ \eta \end{pmatrix}.$$

with domain  $D(L) = \{(u, v, \theta, \eta)^T \in \mathcal{H}_0\}$  such that  $v \in H_2$ ,  $u_{xx} - \theta_{xx} \in H_0$ ,

$$u_{xx} - \int_0^\infty g(s)[c\theta(t-s) - \theta_{xx}(t-s)]ds \in H_0, \quad \eta \in D(T).$$

**Theorem 3.1.** *System (3.1) defines a  $C_0$ -semigroup of contractions  $S(t) = e^{tL}$  on the phase-space  $\mathcal{H}_0$ .*

The proof is done by using the Lumer - Phillips theorem [9, Theorem 4.3].

## 4. LOSS OF EXPONENTIAL STABILITY

To prove the loss of exponential stability we use the following result.

**Theorem 4.1.** *Let  $S(t) = e^{tL}$  be a  $C_0$ -semigroup of contractions in a Hilbert space. Then  $S(t)$  is exponentially stable if and only if,*

$$i\mathbb{R} = \{i\beta : \beta \in \mathbb{R}\} \subset \rho(L) \tag{4.1}$$

and

$$\|(\lambda I - L)^{-1}\| \leq C, \quad \text{for every } \lambda \in i\mathbb{R}. \tag{4.2}$$

The proof of the above theorem can be found in [10] and in [11].

We note that (3.1)-(3.2) is dissipative, because (2.3) implies

$$\langle LU, U \rangle_{\mathcal{H}_0} = \langle T\eta, \eta \rangle_{\mathcal{M}_1} \leq 0, \quad \text{for all } U \in D(L), \tag{4.3}$$

and it is standard matter to show that  $(I - L)$  maps  $D(L)$  onto  $\mathcal{H}_0$ , see [3], where a similar case is treated.

Then, using  $\langle Tu, u \rangle < 0$  for all nonzero  $u$  in  $D(T)$ , one can show that the solution of thermoelastic system (1.1)-(1.2) decays to zero as time approaches  $\infty$ .

Now we are in position of to show our main result.

**Theorem 4.2.** *The semigroup  $S(t) = e^{tL}$  on  $\mathcal{H}_0$  defined by (3.1)-(3.2) is not exponentially stable.*

*Proof.* For  $i\lambda \in \rho(L)$  and  $V = (0, 0, 0, \eta)^T \in \mathcal{H}_0$ , consider the complex equation

$$(i\lambda I - L)U = V \quad (4.4)$$

that when written explicitly reads

$$i\lambda u - v = 0 \quad (4.5)$$

$$i\lambda v - u_{xx} + \theta_{xx} = 0 \quad (4.6)$$

Consider an orthonormal basis  $\{w_j\}_{j \in \mathbb{N}}$  of eigenvectors of the operator  $A$  and the respective eigenvalues  $\{\alpha_n\}_{n \in \mathbb{N}}$ . We recall that  $\alpha_n \rightarrow \infty$  as  $n \rightarrow \infty$ . We set

$$\eta_n(s) = \frac{w_n}{\sqrt{c + \alpha_n}}$$

and

$$V_n = (0, 0, 0, \eta_n)^T.$$

Notice that, using (2.1) and (2.2) we have

$$\|V_n\|_{\mathcal{H}_0} = \|\eta_n\|_{\mathcal{M}_1} = \frac{1}{(c + \alpha_n)}$$

$$\begin{aligned} \int_0^\infty \mu(s) \|B^{1/2} w_n(s)\|^2 ds &= \frac{1}{(c + \alpha_n)} \int_0^\infty \mu(s) (c + \alpha_n) \|w_n(s)\|^2 ds \\ &= \int_0^\infty \mu(s) ds = g(0). \end{aligned}$$

Now we build a sequence of  $\lambda_n$  such that the corresponding solution  $U_n$  of

$$(i\lambda_n I - L)U_n = V_n \quad (4.7)$$

satisfies  $\|U_n\|_{\mathcal{H}_0} \rightarrow \infty$  as  $n \rightarrow \infty$ . In this direction we look for a solution  $U_n = (w_n, w_n, s_n w_n, w_n)$  where  $s_n \in \mathbb{C}$ . Then, from (4.5) and (4.6) we have

$$-\lambda_n^2 - \alpha_n + s_n \alpha_n = 0 \quad (4.8)$$

that implies

$$s_n = 1 + \frac{\lambda_n^2}{\alpha_n}.$$

Choosing  $\lambda_n = |\alpha_n|$  we finally have

$$\|U_n\|_{\mathcal{H}_0} \geq \|s_n w_n\|_{H_0} = |s_n| \geq \frac{\lambda_n^2}{|\alpha_n|} = |\alpha_n| \rightarrow \infty \quad \text{as } n \rightarrow \infty.$$

which yields the conclusion.  $\square$

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