

ON THE FIRST EIGENVALUE OF THE STEKLOV EIGENVALUE PROBLEM FOR THE INFINITY LAPLACIAN

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ABSTRACT. Let Λ_p^p be the best Sobolev embedding constant of $W^{1,p}(\Omega) \hookrightarrow L^p(\partial\Omega)$, where Ω is a smooth bounded domain in \mathbb{R}^N . We prove that as $p \rightarrow \infty$ the sequence Λ_p converges to a constant independent of the shape and the volume of Ω , namely 1. Moreover, for any sequence of eigenfunctions u_p (associated with Λ_p), normalized by $\|u_p\|_{L^\infty(\partial\Omega)} = 1$, there is a subsequence converging to a limit function u_∞ which satisfies, in the viscosity sense, an ∞ -Laplacian equation with a boundary condition.

1. INTRODUCTION

Let Ω be a bounded domain in \mathbb{R}^N with smooth boundary. The minimum Λ_p^p of the Rayleigh quotient, among all nonzero functions in the Sobolev space $W^{1,p}(\Omega)$,

$$\frac{\int_{\Omega} (|\nabla u|^p + |u|^p) dx}{\int_{\partial\Omega} |u|^p dx}$$

is the first eigenvalue of the problem

$$\begin{aligned} -\Delta_p u + |u|^{p-2}u &= 0, & \text{in } \Omega, \\ |\nabla u|^{p-2} \frac{\partial u}{\partial \nu} &= \lambda |u|^{p-2}u, & \text{on } \partial\Omega. \end{aligned} \tag{1.1}$$

where $\Delta_p u = \operatorname{div}(|\nabla u|^{p-2} \nabla u)$ is the p -Laplacian and $\partial/\partial\nu$ is the outer normal derivative along $\partial\Omega$. The eigenvalue problem (1.1) is understood in the weak sense, i.e, $(u, \lambda) \in W^{1,p}(\Omega) \times \mathbb{R}^+$ is an eigenpair if

$$\int_{\Omega} (|\nabla u|^{p-2} \nabla u \nabla v + |u|^{p-2} uv) dx = \lambda \int_{\partial\Omega} |u|^{p-2} uv ds, \quad \forall v \in W^{1,p}(\Omega).$$

The first eigenvalue $\Lambda_p^p = \lambda_1$ is the best constant of the compact embedding $W^{1,p}(\Omega) \hookrightarrow L^p(\partial\Omega)$ and it satisfies

$$\Lambda_p \|u\|_{L^p(\partial\Omega)} \leq \|u\|_{W^{1,p}(\Omega)}, \quad \forall u \in W^{1,p}(\Omega).$$

2000 *Mathematics Subject Classification.* 35J50, 35J55, 35J60, 35J65, 35P30.

Key words and phrases. Nonlinear elliptic equations; eigenvalue problems; p -Laplacian; nonlinear boundary condition; Steklov problem; viscosity solutions.

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Submitted August 4, 2006. Published September 18, 2006.

In this paper we are interested in finding the limit as $p \rightarrow \infty$ of Λ_p . Alternatively, we want to look at the limit of the minimum, as $p \rightarrow \infty$, of the ratio

$$\frac{(\int_{\Omega} (|\nabla u|^p + |u|^p) dx)^{1/p}}{(\int_{\partial\Omega} |u|^p ds)^{1/p}}.$$

It is easy to see that for any positive numbers a and b ,

$$\lim_{p \rightarrow \infty} (a^p + b^p)^{1/p} = \max\{a, b\}.$$

Thus we anticipate that

$$\lim_{p \rightarrow \infty} \frac{(\int_{\Omega} (|\nabla u|^p + |u|^p) dx)^{1/p}}{(\int_{\partial\Omega} |u|^p dx)^{1/p}} = \frac{\max\{\|\nabla u\|_{L^\infty(\Omega)}, \|u\|_{L^\infty(\Omega)}\}}{\|u\|_{L^\infty(\partial\Omega)}}.$$

However, the minimization problem, with minimum value equal 1,

$$\inf_{u \in W^{1,\infty}(\Omega) \setminus \{0\}} \frac{\max\{\|\nabla u\|_{L^\infty(\Omega)}, \|u\|_{L^\infty(\Omega)}\}}{\|u\|_{L^\infty(\partial\Omega)}}, \quad (1.2)$$

has too many solutions. In fact, given a minimizer, we can modify it on any ball inside the domain to obtain another one. The correct Euler-Lagrange equation turns out to be

$$\begin{aligned} \max\{u - |\nabla u|, -\Delta_\infty u\} &= 0, & \text{in } \Omega, \\ \min\{|\nabla u| - u, \frac{\partial u}{\partial \nu}\} &= 0, & \text{on } \partial\Omega, \end{aligned} \quad (1.3)$$

where the operator

$$\Delta_\infty u = \sum_{i,j=1}^N \frac{\partial u}{\partial x_j} \frac{\partial^2 u}{\partial x_j \partial x_i} \frac{\partial u}{\partial x_i},$$

is called the ∞ -Laplacian.

It is clear that for each $p > 1$ the Sobolev embedding constant Λ_p depends on the shape and the volume of domain Ω . We show that when passing to the limit, Λ_p converges to a constant independent of the domain Ω , namely 1. We can also choose a sequence of first eigenfunctions u_p such that the sequence converges uniformly in $C^\alpha(\bar{\Omega})$ to a function u_∞ that satisfies (1.3) in the viscosity sense. In this case we say $(u_\infty, 1)$ is an eigenpair of (1.3). Our main result is:

Theorem 1.1. *For the first eigenvalue of (1.1) we have*

$$\lim_{p \rightarrow \infty} \Lambda_p^{1/p} = 1.$$

For each $p > 1$, let u_p positive eigenfunction be a positive eigenfunction associated with Λ_p^p such that $\|u_p\|_{L^\infty(\partial\Omega)} = 1$. Then there exists a sequence $p_i \rightarrow \infty$ such that $u_{p_i} \rightarrow u_\infty$ in $C^\alpha(\bar{\Omega})$. The limit u_∞ is a solution of (1.3) in the viscosity sense.

To complete the introduction let us mention some recent work on the subject. In [7, 8] the authors study eigenvalue problem for the ∞ -Laplacian with Dirichlet boundary condition. In [1], Steklov eigenvalues for the ∞ -Laplacian are studied when one considers the limit as $p \rightarrow \infty$ of the eigenvalue problem

$$\begin{aligned} -\Delta_p u &= 0, \\ |\nabla u|^{p-2} \frac{\partial u}{\partial \nu} &= \lambda |u|^{p-2} u, & \text{on } \partial\Omega. \end{aligned} \quad (1.4)$$

It is known that the structure of the spectrum of (1.4) is the same as that of (1.1), see [3] and [12]. However, the first eigenvalue of (1.4) is 1 for any $p > 1$ with corresponding constant eigenfunctions in $W^{1,p}(\Omega)$; thus, theorem 1.1 is trivial for problem (1.4). Some arguments and technicalities used here are adapted from [1, 7, 8].

2. Main Results

We first recall the definition of viscosity solutions. Let

$$\begin{aligned} F : \mathbb{R}^N \times \mathbb{R} \times \mathbb{R}^N \times \mathcal{S}(N) &\rightarrow \mathbb{R}, \\ B : \mathbb{R}^N \times \mathbb{R} \times \mathbb{R}^N &\rightarrow \mathbb{R}, \end{aligned}$$

where $\mathcal{S}(N)$ denotes the set of $N \times N$ symmetric matrices.

Consider the boundary-value problem

$$\begin{aligned} F(x, u, \nabla u, D^2 u) &= 0, \\ B(x, u, \nabla u) &= 0, \quad \text{on } \partial\Omega. \end{aligned} \tag{2.1}$$

Definition 2.1. (i) An upper semicontinuous function u is a *viscosity subsolution* of (2.1) if for every $\phi \in C^2(\overline{\Omega})$ such that $u - \phi$ has a strict maximum at the point $x_0 \in \overline{\Omega}$ with $u(x_0) = \phi(x_0)$, we have

$$\begin{aligned} \min\{F(x_0, \phi(x_0), \nabla\phi(x_0), D^2\phi(x_0)), B(x_0, \phi(x_0), \nabla\phi(x_0))\} &\leq 0, \quad x_0 \in \partial\Omega, \\ F(x_0, \phi(x_0), \nabla\phi(x_0), D^2\phi(x_0)) &\leq 0, \quad x_0 \in \Omega. \end{aligned}$$

(ii) A lower semicontinuous function u is a *viscosity supersolution* of (2.1) if for every $\phi \in C^2(\overline{\Omega})$ such that $u - \phi$ has a strict minimum at the point $x_0 \in \overline{\Omega}$ with $u(x_0) = \phi(x_0)$, we have

$$\begin{aligned} \max\{F(x_0, \phi(x_0), \nabla\phi(x_0), D^2\phi(x_0)), B(x_0, \phi(x_0), \nabla\phi(x_0))\} &\geq 0, \quad x_0 \in \partial\Omega, \\ F(x_0, \phi(x_0), \nabla\phi(x_0), D^2\phi(x_0)) &\geq 0, \quad x_0 \in \Omega. \end{aligned}$$

(iii) u is a *viscosity solution* of (2.1) if it is both a supersolution and a subsolution.

In (i) and (ii) the extrema at x_0 need not be strict. We refer to [4] for the theory of viscosity solutions in general and [2] for viscosity solutions with general boundary conditions.

If u is a smooth eigenfunction of (1.1) then by differentiation we get

$$\begin{aligned} -|\nabla u|^{p-2} \Delta u - (p-2)|\nabla u|^{p-4} \Delta_\infty u + |u|^{p-2} u &= 0, \quad \text{in } \Omega, \\ |\nabla u|^{p-2} \frac{\partial u}{\partial \nu} - \lambda |u|^{p-2} u &= 0, \quad \text{on } \partial\Omega. \end{aligned} \tag{2.2}$$

In this case,

$$F(x, z, X, S) = -|X|^{p-2} \text{trace}(S) - (p-2)|X|^{p-4} \langle S \cdot X, X \rangle + |z|^{p-2} z, \tag{2.3}$$

$$B(x, z, X) = |X|^{p-2} \langle X, \nu(x) \rangle - \lambda |z|^{p-2} z. \tag{2.4}$$

It is known that eigenfunctions of (1.1) are in $C^{1,\alpha}(\overline{\Omega})$, see [9, 10] and references therein. Thus it makes sense to talk about viscosity solutions. The following lemma tells us that an eigenfunction is a viscosity solution.

Lemma 2.2. *A weak solution u of (1.1) is a viscosity solution of (2.2).*

Proof. We present the details for the case of supersolutions. Let $x_0 \in \Omega$ and a function $\phi \in C^2(\Omega)$ such that $u(x_0) = \phi(x_0)$ and $u(x) > \phi(x)$, for $x \neq x_0$. We want to show that

$$-|\nabla\phi|^{p-2}\Delta\phi(x_0) - (p-2)|\nabla\phi|^{p-4}\Delta_\infty\phi(x_0) + |u|^{p-2}u(x_0) \geq 0.$$

Suppose that this is not the case, then by continuity there exists a radius $r > 0$ such that, for any $x \in B(x_0, r)$,

$$-|\nabla\phi|^{p-2}\Delta\phi(x) - (p-2)|\nabla\phi|^{p-4}\Delta_\infty\phi(x) + |u|^{p-2}u(x) < 0.$$

Set $m = \inf\{u(x) - \phi(x) : |x - x_0| = r\} > 0$ and let $\Phi(x) = \phi + \frac{1}{2}m$. The function Φ satisfies $\Phi < u$ on $\partial B(x_0, r)$, $\Phi(x_0) > u(x_0)$ and

$$-\operatorname{div}(|\nabla\Phi|^{p-2}\nabla\Phi(x)) + |u|^{p-2}u(x) < 0.$$

Multiplying by $(\Phi - u)^+$ extended by zero outside $B(x_0, r)$ we get

$$\int_{\{\Phi > u\}} |\nabla\Phi|^{p-2}\nabla\Phi \cdot \nabla(\Phi - u) + \int_{\{\Phi > u\}} |u|^{p-2}u(\Phi - u) < 0. \quad (2.5)$$

Since u is a weak solution, we have

$$\int_{\{\Phi > u\}} |\nabla u|^{p-2}\nabla u \cdot \nabla(\Phi - u) + \int_{\{\Phi > u\}} |u|^{p-2}u(\Phi - u) = 0. \quad (2.6)$$

Subtracting (2.6) from (2.5) we get

$$\int_{\{\Phi > u\}} \langle |\nabla\Phi|^{p-2}\nabla\Phi - |\nabla u|^{p-2}\nabla u, \nabla(\Phi - u) \rangle < 0.$$

We obtain a contradiction since the left hand side is bounded below by

$$C(N, p) \int_{\{\Phi > u\}} |\nabla\Phi - \nabla u|^p,$$

where $C(N, p)$ is a positive constant depending only on N and p .

Let λ be the eigenvalue corresponding to u in (1.1). If $x_0 \in \partial\Omega$ and ϕ is a function in $C^2(\overline{\Omega})$ such that $u(x_0) = \phi(x_0)$ and $u(x) > \phi(x)$, for $x \neq x_0$. We want to prove that either

$$-|\nabla\phi|^{p-2}\Delta\phi(x_0) - (p-2)|\nabla\phi|^{p-4}\Delta_\infty\phi(x_0) + |u|^{p-2}u(x_0) \geq 0,$$

$$\text{or } |\nabla\phi|^{p-2}\frac{\partial\phi}{\partial\nu}(x_0) - \lambda|u|^{p-2}u(x_0) \geq 0.$$

Suppose that this is not the case. We repeat the previous argument to obtain

$$\begin{aligned} & \int_{\{\Phi > u\}} |\nabla\Phi|^{p-2}\nabla\Phi \cdot \nabla(\Phi - u) + \int_{\{\Phi > u\}} |u|^{p-2}u(\Phi - u) \\ & < \lambda \int_{\partial\Omega \cap \{\Phi > u\}} |u|^{p-2}u(\Phi - u), \end{aligned}$$

and

$$\begin{aligned} & \int_{\{\Phi > u\}} |\nabla u|^{p-2}\nabla u \cdot \nabla(\Phi - u) + \int_{\{\Phi > u\}} |u|^{p-2}u(\Phi - u) \\ & = \lambda \int_{\partial\Omega \cap \{\Phi > u\}} |u|^{p-2}u(\Phi - u), \end{aligned}$$

which implies

$$\int_{\{\Phi > u\}} \langle |\nabla \Phi|^{p-2} \nabla \Phi - |\nabla u|^{p-2} \nabla u, \nabla(\Phi - u) \rangle < 0.$$

A contradiction is established. This proves that u is a viscosity supersolution. As mentioned, the proof that u is a viscosity subsolution is similar. \square

We are ready to pass to limit as $p \rightarrow \infty$ in the eigenvalue problem. Using the characterization

$$\Lambda_p = \min_{u \neq 0} \frac{(\int_{\Omega} (|\nabla u|^p + |u|^p) dx)^{1/p}}{(\int_{\partial\Omega} |u|^p ds)^{1/p}}, \quad (2.7)$$

one can show the following statement.

Lemma 2.3.

$$\limsup_{p \rightarrow \infty} \Lambda_p \leq 1.$$

Proof. Let $v(x) \equiv 1$, $x \in \bar{\Omega}$. It follows from (2.7) that

$$\Lambda_p \leq \frac{(\int_{\Omega} (|\nabla v|^p + |v|^p) dx)^{1/p}}{(\int_{\partial\Omega} |v|^p ds)^{1/p}} = \frac{|\Omega|^{1/p}}{|\partial\Omega|^{1/p}},$$

where $|\Omega|$ is the Lebesgue measure of Ω and $|\partial\Omega|$ is the boundary measure of $\partial\Omega$. Letting $p \rightarrow \infty$ we obtain the inequality. \square

We recall from [9, 11] that the first eigenvalue Λ_p^p is isolated and simple. Any eigenfunction associated with Λ_p^p is either positive or negative in Ω and any other eigenfunction (not associated with Λ_p^p) has to change sign. We show that Λ_p converges to $\Lambda_{\infty} = 1$, the minimum value of (1.2). We also construct a minimizer of (1.2).

Proposition 2.4. *Given u_p , a positive eigenfunction of (1.1) associated with eigenvalue Λ_p^p , normalized by $\|u_p\|_{L^{\infty}(\partial\Omega)} = 1$. Then there exists a sequence $p_i \rightarrow \infty$ such that $u_{p_i} \rightarrow u_{\infty}$ in $C^{\alpha}(\bar{\Omega})$, where the limit u_{∞} is a minimizer of (1.2) and*

$$\lim_{p \rightarrow \infty} \Lambda_p = 1.$$

Proof. Fix $q > N$. For any $p > q$, one has

$$\begin{aligned} \left(\int_{\Omega} |\nabla u_p|^q + |u_p|^q \right)^{1/q} &\leq |\Omega|^{(1/q)-(1/p)} \left[\int_{\Omega} (|\nabla u_p|^q + |u_p|^q)^{p/q} \right]^{1/p} \\ &\leq (2|\Omega|)^{(1/q)-(1/p)} \left(\int_{\Omega} |\nabla u_p|^p + |u_p|^p \right)^{1/p} \\ &= (2|\Omega|)^{(1/q)-(1/p)} \Lambda_p \left(\int_{\partial\Omega} |u_p|^p \right)^{1/p} \\ &\leq \Lambda_p (2|\Omega|)^{(1/q)-(1/p)} |\partial\Omega|^{1/p}. \end{aligned} \quad (2.8)$$

In above expression, the first inequality follows from a Hölder inequality. We have used that $(a+b)^r \leq 2^{r-1}(a^r + b^r)$, $r \geq 1$, for the second inequality and that $\|u_p\|_{L^{\infty}(\partial\Omega)} = 1$ for the last inequality. We obtain from (2.8) that $\{u_p\}$ is uniformly bounded in $W^{1,q}(\Omega)$. Thus there exists a subsequence $\{u_{p_i}\}$ converging to a function u_{∞} weakly in $W^{1,q}(\Omega)$. Since $q > N$, the Sobolev compact embedding

$W^{1,q}(\Omega) \hookrightarrow C^\alpha(\bar{\Omega})$ holds for any $\alpha \in (0, 1 - N/q)$. It follows that $\{u_{p_i}\}$ converges to u_∞ uniformly in $C^\alpha(\bar{\Omega})$. Moreover, as $p_i \rightarrow \infty$, (2.8) becomes

$$\left(\int_{\Omega} |\nabla u_\infty|^q + |u_\infty|^q \right)^{1/q} \leq \liminf_{p_i \rightarrow \infty} \Lambda_{p_i} (2|\Omega|)^{1/q} \leq (2|\Omega|)^{1/q}. \quad (2.9)$$

On the other hand,

$$\max\{\|\nabla u_\infty\|_{L^q(\Omega)}, \|u_\infty\|_{L^q(\Omega)}\} \leq \left(\int_{\Omega} |\nabla u_\infty|^q + |u_\infty|^q \right)^{1/q} \quad (2.10)$$

Letting $q \rightarrow \infty$, (2.9) and (2.10) imply that

$$\max\{\|\nabla u_\infty\|_{L^\infty(\Omega)}, \|u_\infty\|_{L^\infty(\Omega)}\} \leq 1.$$

The uniform convergence of $\{u_{p_i}\}$ in $C^\alpha(\bar{\Omega})$ gives $\|u_\infty\|_{L^\infty(\partial\Omega)} = 1$. Hence

$$\|\nabla u_\infty\|_{L^\infty(\Omega)} \leq 1 \text{ and } \|u_\infty\|_{L^\infty(\Omega)} = \|u_\infty\|_{L^\infty(\partial\Omega)} = 1.$$

Clearly, u_∞ is the minimizer of (1.2). Furthermore,

$$1 = \|u_\infty\|_{L^\infty(\Omega)} \leq \lim_{q \rightarrow \infty} \left(\int_{\Omega} |\nabla u_\infty|^q + |u_\infty|^q \right)^{1/q} \leq \liminf_{p_i \rightarrow \infty} \Lambda_{p_i},$$

which together with the lemma 2.3 gives $\lim_{p_i \rightarrow \infty} \Lambda_{p_i} = 1$. Since the limit holds for any subsequence, we conclude that $\lim_{p \rightarrow \infty} \Lambda_p = 1$. The proof is complete. \square

Let us verify that the limit of (2.2) as $p \rightarrow \infty$ is (1.3) in the viscosity sense. We obtain from proposition 2.4 that there is a sequence of positive eigenfunctions $\{u_{p_i}\}$ converging to u_∞ uniformly in $\bar{\Omega}$ as $p_i \rightarrow \infty$. Consequently, $u_\infty \geq 0$ in $\bar{\Omega}$.

Lemma 2.5. u_∞ is a viscosity solution of (1.3), i.e.,

$$\begin{aligned} \max\{u_\infty - |\nabla u_\infty|, -\Delta_\infty u_\infty\} &= 0, \quad \text{in } \Omega, \\ \min\{|\nabla u_\infty| - u, \frac{\partial u_\infty}{\partial \nu}\} &= 0, \quad \text{on } \partial\Omega. \end{aligned} \quad (2.11)$$

Proof. First let us check

$$\max\{u_\infty - |\nabla u_\infty|, -\Delta_\infty u_\infty\} = 0 \quad \text{in } \Omega. \quad (2.12)$$

Fix $x_0 \in \Omega$ and a function $\phi \in C^2(\Omega)$ such that $u_\infty(x_0) = \phi(x_0)$ and $u(x) < \phi(x)$, for $x \neq x_0$. Also fix $R > 0$ such that $B(x_0, 2R) \subset \Omega$. For $0 < r < R$ we have

$$\sup\{u_\infty(x) - \phi(x) : x \in B(x_0, R) \setminus B(x_0, r)\} < 0.$$

As $u_{p_i} \rightarrow u_\infty$ uniformly in $\overline{B(x_0, R)}$, for i large enough we conclude that

$$\sup\{u_{p_i}(x) - \phi(x) : x \in B(x_0, R) \setminus B(x_0, r)\} < u_{p_i}(x_0) - \phi(x_0).$$

Therefore for such indices i , $u_{p_i} - \phi$ attains its maximum at $x_i \in B(x_0, r)$. By letting $r \rightarrow 0$ we obtain $x_i \rightarrow x_0$ as $i \rightarrow \infty$. We relabel and denote by $\{x_i\}$ and $\{p_i\}$ the subsequences $\{x_{i_r}\}$ and $\{p_{i_r}\}$. Since u_{p_i} is a subsolution of (2.2) and $x_0 \in \Omega$,

$$-|\nabla \phi|^{p_i-2} \Delta \phi(x_i) - (p_i - 2) |\nabla \phi|^{p_i-4} \Delta_\infty \phi(x_i) + |u_{p_i}|^{p_i-2} u_{p_i}(x_i) \leq 0. \quad (2.13)$$

• Case 1: $\phi(x_0) = u(x_0) > 0$. Then $u_{p_i}(x_i) > 0$ for large i , which implies that $|\nabla \phi(x_i)| \neq 0$ due to (2.13). Dividing by $(p_i - 2) |\nabla \phi(x_i)|^{p_i-4}$ we get

$$-\frac{|\nabla \phi|^2 \Delta \phi(x_i)}{p_i - 2} - \Delta_\infty \phi(x_i) \leq -\left(\frac{u_{p_i}(x_i)}{|\nabla \phi(x_i)|} \right)^{p_i-4} \frac{u_{p_i}^3(x_i)}{p_i - 2}. \quad (2.14)$$

Letting $p_i \rightarrow \infty$ we obtain from (2.14) that

$$\frac{\phi(x_0)}{|\nabla\phi(x_0)|} \leq 1 \text{ and } -\Delta_\infty\phi(x_i) \leq 0.$$

Therefore,

$$\max\{\phi(x_0) - |\nabla\phi(x_0)|, -\Delta_\infty\phi(x_0)\} \leq 0. \tag{2.15}$$

• Case 2: $\phi(x_0) = u(x_0) = 0$. If $|\nabla\phi(x_0)| = 0$, then $\Delta_\infty\phi(x_0) = 0$ and thus (2.15) holds. If $|\nabla\phi(x_0)| \neq 0$, then $|\nabla\phi(x_i)| \neq 0$ for i large. We then obtain (2.14). The right-hand side of (2.14) tends to zero as $p_i \rightarrow \infty$, since

$$\lim_{p_i \rightarrow \infty} \left(\frac{u_{p_i}(x_i)}{|\nabla\phi(x_i)|} \right)^{p_i-4} = 0.$$

Thus $-\Delta_\infty\phi(x_i) \leq 0$ and (2.15) holds in this case. From both cases we conclude that u_∞ is a viscosity subsolution of (2.12).

Next we claim that u_∞ is a viscosity supersolution of (2.12) in Ω . Fix a point $x_0 \in \Omega$ and a function $\phi \in C^2(\Omega)$ such that $u_\infty(x_0) = \phi(x_0)$ and $u_\infty(x) > \phi(x)$, for $x \neq x_0$. We will show that

$$\max\{\phi(x_0) - |\nabla\phi(x_0)|, -\Delta_\infty\phi(x_0)\} \geq 0. \tag{2.16}$$

If $|\nabla\phi(x_0)| = 0$, there is nothing to prove. It suffices to show that if $|\nabla\phi(x_0)| \neq 0$ and $\phi(x_0) - |\nabla\phi(x_0)| < 0$, then $-\Delta_\infty\phi(x_0) \geq 0$. We follow the arguments made in the subsolution case. An analogue of (2.14) is

$$-\frac{|\nabla\phi|^2\Delta\phi(x_i)}{p_i - 2} - \Delta_\infty\phi(x_i) \geq -\left(\frac{u_{p_i}(x_i)}{|\nabla\phi(x_i)|}\right)^{p_i-4} \frac{u_{p_i}^3(x_i)}{p_i - 2}. \tag{2.17}$$

Since $\phi(x_0) - |\nabla\phi(x_0)| < 0$, $\frac{\phi(x_0)}{|\nabla\phi(x_0)|} \leq 1$. Letting $p_i \rightarrow \infty$ it follows from (2.17) that $-\Delta_\infty\phi(x_0) \geq 0$ as claimed. Therefore u_∞ is a viscosity solution of (2.12).

We next need to check on the boundary using definition 2.1. Fix $x_0 \in \partial\Omega$ and a function $\phi \in C^2(\bar{\Omega})$ such that $u_\infty(x_0) = \phi(x_0)$ and $u_\infty(x) < \phi(x)$, for $x \neq x_0$. Using the uniform convergence of u_{p_i} to u_∞ we obtain that $u_{p_i} - \phi$ attains a maximum at $x_i \in \bar{\Omega}$ with $x_i \rightarrow x_0$. If (2.13) holds for infinitely many x_i , we use the argument before to obtain (2.15). Thus we may assume that, for infinitely many $x_i \in \partial\Omega$,

$$|\nabla\phi(x_i)|^{p_i-2} \frac{\partial\phi}{\partial\nu}(x_i) \leq \Lambda_{p_i}^{p_i} |u_{p_i}|^{p_i-2} u_{p_i}(x_i).$$

If $|\nabla\phi(x_0)| = 0$, then $\frac{\partial\phi}{\partial\nu}(x_0) = 0$. If $|\nabla\phi(x_0)| \neq 0$ we obtain

$$\frac{\partial\phi}{\partial\nu}(x_i) \leq \left(\frac{\Lambda_{p_i} |u_{p_i}(x_i)|}{|\nabla\phi(x_i)|}\right)^{p_i-2} \Lambda_{p_i}^2 u_{p_i}(x_i).$$

Since $\Lambda_{p_i} \rightarrow 1$ as $p_i \rightarrow \infty$, we conclude that either

$$\frac{\phi(x_0)}{|\nabla\phi(x_0)|} \geq 1 \text{ or } \frac{\partial\phi}{\partial\nu}(x_0) \leq 0,$$

which implies that $\min\{|\nabla\phi(x_0)| - \phi(x_0), \frac{\partial\phi}{\partial\nu}(x_0)\} \leq 0$. Therefore, at x_0 ,

$$\min \left\{ \max\{\phi - |\nabla\phi|, -\Delta_\infty\phi\}, \min\{|\nabla\phi| - \phi, \frac{\partial\phi}{\partial\nu}\} \right\} \leq 0. \tag{2.18}$$

Fix $x_0 \in \partial\Omega$ and a function $\phi \in C^2(\bar{\Omega})$ such that $u_\infty(x_0) = \phi(x_0)$ and $u_\infty(x) > \phi(x)$, for $x \neq x_0$. Using the uniform convergence of u_{p_i} to u_∞ we obtain that $u_{p_i} - \phi$ attains a minimum at $x_i \in \bar{\Omega}$ with $x_i \rightarrow x_0$. If

$$-|\nabla\phi|^{p_i-2}\Delta\phi(x_i) - (p_i - 2)|\nabla\phi|^{p_i-4}\Delta_\infty\phi(x_i) + |u_{p_i}|^{p_i-2}u_{p_i}(x_i) \geq 0$$

holds for infinitely many x_i , we use the argument before to obtain (2.16). Thus we may assume that, for infinitely many $x_i \in \partial\Omega$,

$$|\nabla\phi(x_i)|^{p_i-2}\frac{\partial\phi}{\partial\nu}(x_i) \geq \Lambda_{p_i}^{p_i}|u_{p_i}|^{p_i-2}u_{p_i}(x_i).$$

If $|\nabla\phi(x_0)| = 0$,

$$\max\{\phi(x_0) - |\nabla\phi|(x_0), -\Delta_\infty\phi(x_0)\} \geq 0$$

If $|\nabla\phi(x_0)| \neq 0$ we obtain

$$\frac{\partial\phi}{\partial\nu}(x_i) \geq \left(\frac{\Lambda_{p_i}|u_{p_i}(x_i)|}{|\nabla\phi(x_i)|}\right)^{p_i-2}\Lambda_{p_i}^2 u_{p_i}(x_i).$$

We conclude that

$$\frac{\phi(x_0)}{|\nabla\phi(x_0)|} \geq 1 \quad \text{and} \quad \frac{\partial\phi}{\partial\nu}(x_0) \geq 0,$$

which implies that $\min\{|\nabla\phi(x_0)| - \phi(x_0), \frac{\partial\phi}{\partial\nu}(x_0)\} \geq 0$. Therefore, at x_0 ,

$$\max\left\{\max\{\phi - |\nabla\phi|, -\Delta_\infty\phi\}, \min\{|\nabla\phi| - \phi, \frac{\partial\phi}{\partial\nu}\}\right\} \geq 0. \quad (2.19)$$

Inequalities (2.18) and (2.19) prove that u_∞ satisfies in the viscosity sense the boundary condition of (1.3). \square

Theorem 1.1 follows from proposition 2.4 and lemma 2.5.

Acknowledgments. The author is grateful to Professor Juan Manfredi for many interesting discussions which leads to the current work. The author also would like to thank Professor Klaus Schmitt for his careful reading of the manuscript and his valuable comments. This work was done while the author was visiting the MSRI in Berkeley, California, in the fall of 2005 and Utah State University in the spring of 2006.

REFERENCES

- [1] J. Garcia-Azorero, J. J. Manfredi, I. Peral and J. D. Rossi; *Steklov eigenvalues for the ∞ -Laplacian*, manuscript.
- [2] G. Barles, *Fully nonlinear Neumann type boundary conditions for second-order elliptic and parabolic equations*, J. Differential Equations, 106 (1993), pp. 90–106.
- [3] J. F. Bonder and J. D. Rossi, *Existence results for the p -Laplacian with nonlinear boundary conditions*, J. Math. Anal. Appl., 263 (2001), pp. 195–223.
- [4] M. G. Crandall, H. Ishii, and P.-L. Lions, *User's guide to viscosity solutions of second order partial differential equations*, Bull. Amer. Math. Soc. (N.S.), 27 (1992), pp. 1–67.
- [5] L. C. Evans, *Partial Differential Equations*, American Mathematical Society, Providence, RI, 1998.
- [6] D. Gilbarg and N. S. Trudinger, *Elliptic Partial Differential Equations of Second Order*, Classics in Mathematics, Springer-Verlag, Berlin, 2001. Reprint of the 1998 edition.
- [7] P. Juutinen and P. Lindqvist, *On the higher eigenvalues for the ∞ -eigenvalue problem*, Calc. Var. Partial Differential Equations, 23 (2005), pp. 169–192.
- [8] P. Juutinen, P. Lindqvist, and J. J. Manfredi, *The ∞ -eigenvalue problem*, Arch. Ration. Mech. Anal., 148 (1999), pp. 89–105.
- [9] A. Lê, *Eigenvalue problems for the p -Laplacian*, Nonlinear Analysis, TMA, 64 (2006), pp. 1057–1099.

- [10] G. M. Lieberman, *Boundary regularity for solutions of degenerate elliptic equations*, Nonlinear Anal., 12 (1988), pp. 1203–1219.
- [11] S. Martínez and J. D. Rossi, *Isolation and simplicity for the first eigenvalue of the p -Laplacian with a nonlinear boundary condition*, Abstr. Appl. Anal., 7 (2002), pp. 287–293.
- [12] O. Torné, *Steklov problem with an indefinite weight for the p -Laplacian*, Electron. J. Differential Equations, 2005 (2005), No. 87, 8 pp. (electronic).

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