

## TRANSPORT OPERATOR ON PHASE SPACES WITH FINITE TIME OF SOJOURN PROPERTY

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ABSTRACT. In this article, the transport operator with general boundary conditions is discussed. According to a smallness hypothesis on the boundary operator and to finite time of sojourn property of phase spaces, we prove that the transport operator generates a strongly continuous semigroup and we give its upper bound.

### 1. INTRODUCTION

This article concerns the transport equation

$$\frac{\partial f}{\partial t}(x, v) = -v \cdot \nabla_x f(x, v), \quad (x, v) \in \Omega \quad (1.1)$$

where,  $\Omega = X \times V$  with  $X \subset \mathbb{R}^n$  ( $n \geq 1$ ) is a bounded open subset with smooth boundary  $\partial X$  and  $d\mu$  is a Radon measure on  $\mathbb{R}^n$  with bounded support  $V$ . If we denote by  $\Gamma_-$  (resp.  $\Gamma_+$ ) the incoming (resp. outgoing) part of the phase space boundary  $\Gamma = \partial X \times V$ , then the boundary condition is modelled as

$$f(t)|_{\Gamma_-} = K(f(t)|_{\Gamma_+}) \quad (1.2)$$

where,  $f(t)|_{\Gamma_-}$  (resp.  $f(t)|_{\Gamma_+}$ ) is the incoming (resp. outgoing) particle flux. The boundary operator  $K$  is bounded linear into suitable function spaces on  $\Gamma_-$  and  $\Gamma_+$  (for more explanations see next Section). All known boundary conditions (vacuum, specular reflections, periodic, ...) are special examples of our general context.

If  $\|K\| \leq 1$ , it is well known, in the pioneer works [1, 9, 10], that the transport model (1.1)–(1.2) is governed by a strongly continuous semigroup of contractions. However, the case  $\|K\| > 1$  has been rarely studied and some contributions are made in [2, 3, 5, 6]. There is another contribution made in [8] (see last Section).

The difficulty concerning the case  $\|K\| > 1$  is closely related to the increasing number of, on one hand, the incoming particles whose the time of sojourn  $\tau(x, v)$  may be arbitrary small and on the other hand, to the particles in  $X$  of which the time of sojourn  $t(x, v)$  may be arbitrary big. In order to take into account such as particles, we intuitively have to set hypotheses on the geometry of  $(X, V)$  and on the boundary operators  $K$ . So, the first hypothesis concerns boundary operators  $K$  satisfying

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(H1) There exists  $\varepsilon_0 > 0$  such that  $\|K\chi_{\varepsilon_0}\|_{\mathcal{L}(L^1(\Gamma_+), L^1(\Gamma_-))} < 1$ , where the characteristic operator  $\chi_{\varepsilon_0} \in \mathcal{L}(L^1(\Gamma_+))$  is defined by

$$\chi_{\varepsilon_0}\psi(x, v) = \begin{cases} \psi(x, v) & \text{if } \tau(x, v) \leq \varepsilon_0 \\ 0 & \text{otherwise.} \end{cases} \quad (1.3)$$

The second hypothesis acts on the geometry of  $V$  in the following sense

(H2)  $0 \notin V$

which intuitively leads us to set our definition

**Definition 1.1.** The phase space  $(X, V)$  has *finite time of sojourn property* if

$$T_{\max} := \sup_{(x, v) \in \Omega} t(x, v) < \infty. \quad (1.4)$$

Clearly, the hypothesis (H2) implies that velocities cannot vanish; therefore Definition 1.1 holds because of the boundedness of  $X$ . For instance, let the phase space  $((0, 1) \times (a, b))$  ( $a > 0$ ) be related to a model of cell dynamic populations already studied in [7]. The phase space  $((0, 1) \times (a, b))$  ( $a > 0$ ) fulfils the definition above because of  $T_{\max} = \frac{1}{a} < \infty$ .

In this paper, we discuss the case  $\|K\| > 1$  and at this end, we suppose that the hypotheses (H1) and (H2) hold. So, we prove that the transport model (1.1)–(1.2) is governed by a strongly continuous semigroup and we give its upper bound. We end this paper by remarks and comments.

## 2. SETTING OF THE PROBLEM

In this section we state preparatory Lemmas for the next Section. Let us consider the Banach space  $L^1(\Omega)$  whose natural norm is

$$\|\varphi\|_1 = \int_{\Omega} |\varphi(x, v)| dx d\mu(v) \quad (2.1)$$

where,  $\Omega = X \times V$ . We set  $n(x)$  the outer unit normal at  $x \in \partial X$ , where, the boundary  $\partial X$  is equipped with the Lebesgue measure  $d\gamma$  and we denote

$$\Gamma_{\pm} = \{(x, v) \in \Gamma : \pm v \cdot n(x) > 0\},$$

where  $\Gamma = \partial X \times V$ . For each  $(x, v) \in \Omega$ , we set

$$t(x, v) = \inf\{t : t > 0, x - tv \notin X\},$$

the time of sojourn in  $X$ , and

$$\theta(x, v) = t(x, v) + t(x, -v),$$

the chord of sojourn. Similarly, if  $(x, v) \in \Gamma_+$  we set

$$\tau(x, v) = \inf\{t : t > 0, x - tv \notin X\}.$$

Next, we introduce the partial Sobolev space

$$W^1(\Omega) = \{\varphi \in L^1(\Omega) : v \cdot \nabla_x \varphi \in L^1(\Omega), \theta^{-1}\varphi \in L^1(\Omega)\}$$

whose norm is

$$\|\varphi\|_{W^1(\Omega)} = \|v \cdot \nabla_x \varphi\|_1 + \|\theta^{-1}\varphi\|.$$

Finally, we consider the trace spaces  $L^1(\Gamma_{\pm})$  endowed with the norm

$$\|\varphi\|_{L^1(\Gamma_{\pm})} = \int_{\Gamma_{\pm}} |\varphi(x, v)| d\xi$$

where,  $d\xi = |v \cdot n(x)|d\gamma d\mu(v)$ . In this context, we define the following trace mapping

$$\gamma_+\varphi = \varphi|_{\Gamma_+} \quad \text{and} \quad \gamma_-\varphi = \varphi|_{\Gamma_-}$$

for which we have our new result.

**Lemma 2.1** ([3]). *The trace mappings*

$$\gamma_+ : W^1(\Omega) \rightarrow L^1(\Gamma_+) \quad \text{and} \quad \gamma_- : W^1(\Omega) \rightarrow L^1(\Gamma_-)$$

*are continuous, surjective and admit continuous lifting operators.*

Let  $K$  be a bounded linear operator from  $L^1(\Gamma_+)$  into  $L^1(\Gamma_-)$ . So, it is clear that above Lemma allows us to give a sense to the following transport operator

$$\begin{aligned} T_K\varphi &= -v \cdot \nabla_x \varphi \quad \text{on the domain} \\ D(T_K) &= \{\varphi \in W^1(\Omega) : \gamma_-\varphi = K\gamma_+\varphi\}. \end{aligned} \tag{2.2}$$

If the boundary operator satisfies  $K = 0$ , then the corresponding operator  $T_0$  is defined as follows

$$\begin{aligned} T_0\varphi &= -v \cdot \nabla_x \varphi \quad \text{on the domain} \\ D(T_0) &= \{\varphi \in W^1(\Omega) : \gamma_-\varphi = 0\} \end{aligned} \tag{2.3}$$

has some properties summarized next.

**Lemma 2.2.** *We have*

- (1) *The operator  $T_0$  generates, on  $L^1(\Omega)$ , a strongly continuous semigroup of contractions  $(U_0(t))_{t \geq 0}$ . Furthermore,  $U_0(t)$  is a positive operator; i.e.,  $U_0(t)\varphi \geq 0$  for all  $\varphi \in (L^1(\Omega))_+$ .*
- (2) *Let  $\lambda > 0$  be fixed. Then, for all  $\varphi \in (L^1(\Omega))_+ - \{0\}$  we have  $(\lambda - T_0)^{-1}\varphi \in (L^1(\Omega))_+ - \{0\}$  and  $\gamma_+(\lambda - T_0)^{-1}\varphi \in (L^1(\Gamma_+))_+ - \{0\}$ .*
- (3) *Let  $\lambda > 0$ . Then*

$$\|(\lambda - T_0)^{-1}g\|_1 \leq \frac{\|g\|_1}{\lambda}, \tag{2.4}$$

$$\|\theta^{-1}(\lambda - T_0)^{-1}g\|_1 \leq \|g\|_1 \tag{2.5}$$

for all  $g \in L^1(\Omega)$ .

*Proof.* The items (1), (2) and (3) follow easily from

$$(\lambda - T_0)^{-1}g(x, v) = \int_0^{t(x,v)} e^{-\lambda s} g(x - sv, v) ds$$

where,  $\lambda > 0$  and  $g \in L^1(\Omega)$ . □

**Lemma 2.3.** *Let  $A$  be the operator*

$$A\psi(x, v) = \psi(x - \tau(x, v)v, v). \tag{2.6}$$

*Then  $A$  is a positive isometry from  $L^1(\Gamma_-)$  to  $L^1(\Gamma_+)$ ; i.e.,*

$$\|A\psi\|_{L^1(\Gamma_+)} = \|\psi\|_{L^1(\Gamma_-)} \tag{2.7}$$

for all  $\psi \in L^1(\Gamma_+)$ .

*Proof.* Let  $\psi \in L^1(\Gamma_+)$ . As  $u(x, v) = \psi(x - t(x, v)v, v)$  is the unique solution of the boundary value problem

$$\begin{aligned} v \cdot \nabla_x u &= 0 \\ \gamma_- u &= \psi. \end{aligned}$$

Then multiplying the first equation by  $(\operatorname{sgn} u)$  and using

$$(\operatorname{sgn} u)v \cdot \nabla_x u = v \cdot \nabla_x (|u|), \quad (2.8)$$

we obtain  $v \cdot \nabla_x (|u|) = 0$ . Integrating this equation over  $\Omega$  and using Green's identity, we obtain

$$\int_{\Gamma_+} |\gamma_+ u(x, v)| d\xi = \int_{\Gamma_-} |\gamma_- u(x, v)| d\xi;$$

therefore,

$$\int_{\Gamma_+} |A\psi(x, v)| d\xi = \int_{\Gamma_-} |\psi(x, v)| d\xi$$

whence (2.7). The positivity of  $A$  is obvious.  $\square$

### 3. GENERATION THEOREM

In this section, we are only concerned with boundary operators whose norm satisfies  $\|K\| > 1$ . So, according to (H1)–(H2), we prove that the transport operator  $T_K$  given by (2.2) generates, on  $L^1(\Omega)$ , a strongly continuous semigroup. Before we state this main goal, we have to show the following lemmas.

**Lemma 3.1.** *Let  $K$  be a boundary operator with  $\|K\| > 1$  and suppose that (H1) holds. Let  $K_\lambda$  ( $\lambda \geq 0$ ) be the operator*

$$K_\lambda \psi := K(\alpha_\lambda \psi), \quad (3.1)$$

where

$$\alpha_\lambda(x, v) = e^{-\lambda\tau(x, v)}.$$

Then  $K_\lambda$  is a bounded linear operator from  $L^1(\Gamma_+)$  to  $L^1(\Gamma_-)$ . Furthermore, we have

$$\lambda > \omega_0 \implies \|K_\lambda\| < 1, \quad (3.2)$$

$$\|K_{\omega_0}\| \leq 1, \quad (3.3)$$

where

$$\omega_0 = \frac{1}{\varepsilon_0} \ln \|K\|. \quad (3.4)$$

Moreover, if  $K$  is a positive operator, then  $K_\lambda$  is also a positive operator.

*Proof.* Let  $\lambda \geq 0$ . For all  $\psi \in L^1(\Gamma_+)$  we obviously have  $\chi_{\varepsilon_0}^2 = \chi_{\varepsilon_0}$  which implies

$$\alpha_\lambda \psi = \chi_{\varepsilon_0}^2 (\alpha_\lambda \psi) + \bar{\chi}_{\varepsilon_0} (\alpha_\lambda \psi);$$

therefore,

$$K_\lambda \psi = K\chi_{\varepsilon_0} (\chi_{\varepsilon_0} \alpha_\lambda \psi) + K\bar{\chi}_{\varepsilon_0} (\alpha_\lambda \psi),$$

where the characteristic operator  $\chi_{\varepsilon_0}$  is given by (1.3), and  $\bar{\chi}_{\varepsilon_0}$  is the characteristic operator

$$\bar{\chi}_{\varepsilon_0} \psi(x, v) = \begin{cases} \psi(x, v) & \text{if } \tau(x, v) > \varepsilon_0 \\ 0 & \text{otherwise.} \end{cases}$$

This implies

$$\begin{aligned} \|K_\lambda \psi\|_{L^1(\Gamma_-)} &\leq \|K\chi_{\varepsilon_0}(\chi_{\varepsilon_0}\alpha_\lambda\psi)\|_{L^1(\Gamma_-)} + \|K\bar{\chi}_{\varepsilon_0}(\alpha_\lambda\psi)\|_{L^1(\Gamma_-)} \\ &\leq \|K\chi_{\varepsilon_0}\| \|\chi_{\varepsilon_0}\alpha_\lambda\psi\|_{L^1(\Gamma_+)} + \|K\| \|\bar{\chi}_{\varepsilon_0}(\alpha_\lambda\psi)\|_{L^1(\Gamma_+)} \\ &\leq \|K\chi_{\varepsilon_0}\| \|\chi_{\varepsilon_0}\psi\|_{L^1(\Gamma_+)} + e^{-\lambda\varepsilon_0} \|K\| \|\bar{\chi}_{\varepsilon_0}\psi\|_{L^1(\Gamma_+)} \\ &\leq \max\{\|K\chi_{\varepsilon_0}\|, e^{-\lambda\varepsilon_0}\|K\|\} \{\|\chi_{\varepsilon_0}\psi\| + \|\bar{\chi}_{\varepsilon_0}\psi\|\} \\ &= \{\|K\chi_{\varepsilon_0}\|, e^{-\lambda\varepsilon_0}\|K\|\} \|\psi\|_{L^1(\Gamma_+)} \end{aligned}$$

which leads to

$$\|K_\lambda\| \leq \max\{\|K\chi_{\varepsilon_0}\|, e^{-\lambda\varepsilon_0}\|K\|\}.$$

Now, the above relation clearly implies

$$\|K_\lambda\| < 1 \quad \text{if } \lambda > \omega_0$$

and  $\|K_{\omega_0}\| \leq 1$ . Finally, if  $K$  is a positive operator, the positivity of the operator  $K_\lambda$  is then obvious. The proof is now achieved.  $\square$

Thanks Lemma above, the resolvent operator of (2.2) is given as follows.

**Lemma 3.2.** *Let  $K$  be a boundary operator whose satisfying  $\|K\| > 1$ , and suppose that (H1) holds. Then, for all  $\lambda > \omega_0$ , we have  $\lambda \in \rho(T_K)$  and*

$$\begin{aligned} (\lambda - T_K)^{-1}g(x, v) &= (\lambda - T_0)^{-1}g(x, v) + \\ e^{-\lambda t(x, v)}(I - K_\lambda A)^{-1}K\gamma_-(\lambda - T_0)^{-1}g(x - t(x, v)v, v) \end{aligned} \quad (3.5)$$

for almost all  $(x, v) \in \Omega$  and for all  $g \in L^1(\Omega)$ , where,  $A$  is the operator given by (2.6). Furthermore, if  $K$  is a positive operator,  $(\lambda - T_K)^{-1}$  is then a positive operator for all  $\lambda > \omega_0$ .

*Proof.* Let  $\lambda > \omega_0$ . For all  $g \in L^1(\Omega)$ , the general solution of

$$\lambda\varphi = -v \cdot \nabla_x \varphi + g, \quad (3.6)$$

is given by

$$\varphi(x, v) = e^{-\lambda t(x, v)}\psi(x - t(x, v)v, v) + (\lambda - T_0)^{-1}g(x, v), \quad (3.7)$$

for almost all  $(x, v) \in \Omega$ , where  $T_0$  is already studied in Lemma 2.2 and  $\psi$  is any function of  $L^1(\Gamma_-)$ . In the sequel, let us prove that  $\varphi \in D(T_K)$ .

Integrating (3.7) over  $\Omega$ , a simple calculation together with (2.4) give us

$$\begin{aligned} \|\varphi\|_1 &\leq \int_\Omega e^{-\lambda t(x, v)} |\psi(x - t(x, v)v, v)| dx d\mu(v) + \|(\lambda - T_0)^{-1}g\|_1 \\ &\leq \frac{1}{\lambda} \|\psi\|_{L^1(\Gamma_-)} + \frac{\|g\|_1}{\lambda} < \infty \end{aligned}$$

which implies, by (3.6), that

$$\|v \cdot \nabla_x \varphi\|_1 \leq \lambda \|\varphi\|_1 + \|g\|_1 < \infty.$$

Multiplying (3.7) by  $\theta^{-1}$  and integrating it over  $\Omega$ , a simple calculation together with (2.5) lead to

$$\|\theta^{-1}\varphi\|_1 \leq \|\psi\|_{L^1(\Gamma_-)} + \|g\|_1 < \infty;$$

therefore,  $\varphi \in W^1(\Omega)$ . Next,  $\varphi$  satisfies  $\gamma_- \varphi = K\gamma_- \varphi$  if and only if  $\psi$  satisfies

$$\psi = K_\lambda A\psi + K\gamma_-(\lambda - T_0)^{-1}g. \quad (3.8)$$

By (2.7) and (3.2) we obtain

$$\|K_\lambda A\| \leq \|K_\lambda\| \|A\| < 1; \quad (3.9)$$

therefore, (3.8) admits the unique solution

$$\psi = (I - K_\lambda A)^{-1} K \gamma_-(\lambda - T_0)^{-1} g$$

which we put in (3.7) to obtain (3.5). In order to achieve the proof, it suffices to show the positivity of the operator  $(\lambda - T_K)^{-1}$ .

Let  $g \in (L^1(\Omega))_+$  and note that the positivity of the operator  $K$  implies that of the operator  $K_\lambda$ . By (3.5) and the second item of Lemma 2.2 we obtain

$$(\lambda - T_K)^{-1} g(x, v) \geq e^{-\lambda t(x, v)} (I - K_\lambda A)^{-1} K \gamma_+(\lambda - T_0)^{-1} g(x - t(x, v)v, v)$$

for almost all  $(x, v) \in \Omega$ . Thanks to (3.9) we have

$$(I - K_\lambda A)^{-1} K = \left( \sum_{n=0}^{\infty} (K_\lambda A)^n \right) K \geq IK = K$$

therefore,

$$(\lambda - T_K)^{-1} g(x, v) \geq e^{-\lambda t(x, v)} K \gamma_+(\lambda - T_0)^{-1} g(x - t(x, v)v, v)$$

for almost all  $(x, v) \in \Omega$ . Finally, the positivity of  $K$  and the second item of Lemma 2.2 clearly imply the positivity of  $(\lambda - T_K)^{-1} g$ . The proof is now achieved.  $\square$

Now, we are ready to state the main result of this work.

**Theorem 3.3.** *Let  $K$  be a boundary operator with  $\|K\| > 1$ , and suppose that (H1)–(H2) hold. Then, the transport operator  $T_K$  given by (2.2) generates, on  $L^1(\Omega)$ , a strongly continuous semigroup  $(U_K(t))_{t \geq 0}$  satisfying*

$$\|U_K(t)g\|_1 \leq e^{\omega_0(T_{\max} + t)} \|g\|_1 \quad t \geq 0, \quad (3.10)$$

for all  $g \in L^1(\Omega)$ , where,  $T_{\max}$  and  $\omega_0$  are given by (1.4) and (3.4). Furthermore, if  $K$  is a positive operator,  $(U_K(t))_{t \geq 0}$  is positive too.

*Proof.* First, let us define on  $L^1(\Omega)$  the norm

$$\|g\|_1 = \int_{\Omega} |g(x, v)| h(x, v) dx d\mu(v) \quad (3.11)$$

where,  $h(x, v) = e^{\omega_0 t(x, v)}$ . By (H2), (1.4) holds; therefore, the norms (2.1) and (3.11) are equivalent because

$$\|g\|_1 \leq \|g\|_1 \leq e^{\omega_0 T_{\max}} \|g\|_1 \quad (3.12)$$

for all  $g \in L^1(\Omega)$ .

Next, let  $\lambda > \omega_0$  and  $g \in L^1(\Omega)$ . Thanks to Lemma 3.2 we obtain that

$$\varphi = (\lambda - T_K)^{-1} g \in D(T_K) \quad (3.13)$$

is the unique solution of  $\lambda\varphi = T_K\varphi + g$ . Therefore  $\varphi$  satisfies

$$\lambda\varphi = -v \cdot \nabla_x \varphi + g, \quad (3.14)$$

$$\gamma_- \varphi = K \gamma_- \varphi. \quad (3.15)$$

Multiplying (3.14) by  $(\text{sgn } \varphi)h$  and integrating it over  $\Omega$ ,

$$\begin{aligned} \lambda \|\varphi\|_1 &= - \int_{\Omega} v \cdot \nabla_x (|\varphi|) h(x, v) dx d\mu(v) + \int_{\Omega} ((\text{sgn } \varphi)hg)(x, v) dx d\mu(v) \\ &:= I + J. \end{aligned} \quad (3.16)$$

Integrating by parts,

$$\begin{aligned} I &= - \int_{\Omega} v \cdot \nabla_x (|h\varphi|)(x, v) \, dx \, d\mu(v) + \omega_0 \int_{\Omega} |(h\varphi)(x, v)| \, dx \, d\mu(v) \\ &= \int_{\Gamma_-} |\gamma_-(h\varphi)(x, v)| \, d\xi - \int_{\Gamma_+} |\gamma_+(h\varphi)(x, v)| \, d\xi + \omega_0 \|\varphi\|_1 \\ &= \int_{\Gamma_-} |\gamma_-\varphi(x, v)| \, d\xi - \int_{\Gamma_+} |\gamma_+(h\varphi)(x, v)| \, d\xi + \omega_0 \|\varphi\|_1. \end{aligned}$$

By (3.15) and the fact that  $(\gamma_+h)\alpha_{\omega_0} = 1$ , we obtain

$$\begin{aligned} I &= \int_{\Gamma_-} |K\gamma_+\varphi(x, v)| \, d\xi - \int_{\Gamma_+} |\gamma_+(h\varphi)(x, v)| \, d\xi + \omega_0 \|\varphi\|_1 \\ &= \int_{\Gamma_-} |K(\alpha_{\omega_0}\gamma_+(h\varphi))(x, v)| \, d\xi - \int_{\Gamma_+} |\gamma_+(h\varphi)(x, v)| \, d\xi + \omega_0 \|\varphi\|_1 \\ &= \int_{\Gamma_-} |K_{\omega_0}(\gamma_+(h\varphi))(x, v)| \, d\xi - \int_{\Gamma_+} |\gamma_+(h\varphi)(x, v)| \, d\xi + \omega_0 \|\varphi\|_1 \\ &\leq (\|K_{\omega_0}\| - 1) \|\gamma_+(h\varphi)\|_{L^1(\Gamma_+)} + \omega_0 \|\varphi\|_1; \end{aligned}$$

therefore

$$I \leq \omega_0 \|\varphi\|_1 \tag{3.17}$$

because of (3.3). For the term  $J$ , we obviously have

$$J = \int_{\Omega} ((\text{sgn } \varphi)hg)(x, v) \, dx \, d\mu(v) \leq \|g\|_1. \tag{3.18}$$

Putting now (3.13), (3.17) and (3.18) in (3.16) we obtain

$$\|(\lambda - T_K)^{-1}g\|_1 \leq \frac{\|g\|_1}{(\lambda - \omega_0)}.$$

Thanks to Hille-Yosida's theorem, the operator  $T_K$  generates on  $L^1(\Omega)$  a strongly continuous semigroup  $(U_K(t))_{t \geq 0}$  satisfying

$$\|U_K(t)g\|_1 \leq e^{t\omega_0} \|g\|_1 \quad t \geq 0 \tag{3.19}$$

for all  $g \in L^1(\Omega)$ . Now (3.10) follows from (3.12) and (3.19). In order to achieve the proof, it suffices to show the positivity of the semigroup  $(U_K(t))_{t \geq 0}$ .

Let  $g \in (L^1(X \times V))_+$  and  $t > 0$ . Lemma 3.2 leads to

$$\left(\frac{n}{t} - T_K\right)^{-1}g \geq 0$$

for  $n$  large enough. Now, the exponential formula

$$U_K(t)g = \lim_{n \rightarrow \infty} \left[\frac{n}{t} \left(\frac{n}{t} - T_K\right)^{-1}\right]^n g \geq 0$$

achieves the proof. □

We finish this section by giving an example of a boundary operator  $K$  satisfying our hypothesis (H1).

**Lemma 3.4.** *Let  $K \in \mathcal{L}(L^1(\Gamma_+), L^1(\Gamma_-))$  be a Maxwell boundary operator; i.e.,  $K = C + B$  where,*

$$C\psi(x, v) = \int_{\Gamma_+} k(x, v, x', v') \psi(x', v') |v' \cdot n(x')| \, d\gamma(x') \, d\mu(v') \quad (x, v) \in \Gamma_-$$

with  $k \geq 0$  and  $B \in \mathcal{L}(L^1(\Gamma_+), L^1(\Gamma_-))$  is a given operator such that  $\|B\| < 1$ . If

$$\limsup_{\varepsilon \rightarrow 0} \operatorname{ess\,sup}_{\{\tau(y, v') \leq \varepsilon\}} \int_{\Gamma_-} k(x, v, y, v') |v \cdot n(x)| d\gamma(x) d\mu(v') < 1 - \|B\|,$$

then the hypothesis (H1) holds.

*Proof.* It is clear that there exists  $\varepsilon_0 > 0$  such that

$$\|C\chi_{\varepsilon_0}\| = \operatorname{ess\,sup}_{\{\tau(y, v') \leq \varepsilon_0\}} \int_{\Gamma_-} k(x, v, y, v') |v \cdot n(x)| d\gamma(x) d\mu(v') < 1 - \|B\|;$$

therefore

$$\|K\chi_{\varepsilon_0}\| \leq \|C\chi_{\varepsilon_0}\| + \|B\chi_{\varepsilon_0}\| < 1 - \|B\| + \|B\| = 1.$$

The proof is achieved.  $\square$

#### 4. REMARKS AND COMMENTS

As we pointed in the introduction, this section deals with some remarks and comments on [8], using our notation.

**Remark 4.1.** In [8, page 288, line 14], the authors claim that the traces  $\psi|_{\Gamma_{\pm}}$  ( $= \gamma_{\pm}\psi$ ) of  $\psi$  in

$$\mathcal{W}^1(\Omega) = \{\varphi \in L^1(\Omega), v \cdot \nabla_x \varphi \in L^1(\Omega)\},$$

are well defined and belong to  $L^1(\Gamma_{\pm})$ . According to our Lemma 2.1, this claim is incorrect.

**Remark 4.2.** Note that [8, Theorem 5.2] is incorrect. Indeed, the authors consider positive boundary operators  $K$  satisfying

(H1)

$$\lim_{\varepsilon \rightarrow 0} \|K\chi_{\varepsilon}\|_{\mathcal{L}(L^1(\Gamma_+), L^1(\Gamma_-))} < 1$$

where the characteristic operator  $\chi_{\varepsilon_0}$  is given by (1.3), and

(H2)

$$\|K\psi\|_{L^1(\Gamma_-)} \geq \|\psi\|_{L^1(\Gamma_+)}$$

for all  $\psi \in (L^1(\Gamma_+))_+$ .

According to (H1)–(H2), the authors claim that the operator  $T_K$  defined by

$$T_K\varphi = -v \cdot \nabla_x \varphi \quad \text{on the domain}$$

$$D(T_K) = \{\varphi \in \mathcal{W}^1(\Omega), \gamma_- \varphi = K\gamma_+ \varphi\}$$

generates, on  $L^1(\Omega)$ , a strongly continuous semigroup.

However, by (H1), there exist  $\varepsilon_0 > 0$  and

$$0 < \alpha < 1 \tag{4.1}$$

such that

$$\|K\chi_{\varepsilon_0}\|_{\mathcal{L}(L^1(\Gamma_+), L^1(\Gamma_-))} < \alpha \tag{4.2}$$

which implies that

$$\|K\chi_{\varepsilon_0}\psi\|_{L^1(\Gamma_-)} < \alpha \|\psi\|_{L^1(\Gamma_+)} \tag{4.3}$$

for all  $\psi \in L^1(\Gamma_+)$ . Next, let us consider  $\psi \in (L^1(\Gamma_+))_+$  such that  $\chi_{\varepsilon_0}\psi \neq 0$ . Now, clearly the fact that  $\chi_{\varepsilon_0}^2 = \chi_{\varepsilon_0}$  together with the hypothesis (H2) and (4.3) lead us to

$$\|\chi_{\varepsilon_0}\psi\| = \|\chi_{\varepsilon_0}^2\psi\| \leq \|K(\chi_{\varepsilon_0}^2\psi)\| = \|K\chi_{\varepsilon_0}(\chi_{\varepsilon_0}\psi)\| < \alpha \|\chi_{\varepsilon_0}\psi\|$$



whence  $1 < \alpha$  which contradicts (4.1). Therefore, there are no positive boundary operators  $K$  satisfying simultaneously  $(\overline{\text{H1}})$  and  $(\overline{\text{H2}})$ .

Finally, note that (4.1) and (4.2) imply

$$\|K\chi_{\varepsilon_0}\|_{\mathcal{L}(L^1(\Gamma_+), L^1(\Gamma_-))} < 1,$$

and our hypothesis (H1) holds.

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