Electronic Journal of Differential Equations, Vol. 2024 (2024), No. 61, pp. 1–15.

ISSN: 1072-6691. URL: https://ejde.math.txstate.edu, https://ejde.math.unt.edu

DOI: 10.58997/ejde.2024.61

EXISTENCE AND FORMS OF ENTIRE SOLUTIONS TO SYSTEM OF NON-LINEAR PARTIAL DIFFERENTIAL EQUATIONS

ABHIJIT BANERJEE, JHUMA SARKAR

ABSTRACT. The main objective of this article is to explore the existence and forms of transcendental entire solutions of some systems of non-linear partial differential equations. We obtain two results and illustrate the results with several examples. This article improves the results in [5, 15]. In the last section we discuss the differences between the solutions involving homogeneous and non-homogeneous operators, and state an open question for the sake of future research.

1. Introduction

The development of the difference analogue of the Nevanlinna theory [4, 11] has greatly influenced the study of difference and difference-differential equations. Naturally, this topic has become a central focus for many researchers in the field. In 1966, Gross [6] studied the existence and form of transcendental entire solution of the equation $f(z)^m + g(z)^m = 1$, and settled the problem for m = 2 and pointed out that the equation does not possess any non-constant transcendental entire solution if m > 2. This significant result opened new avenues for further exploration about the existence and form of transcendental entire solutions for variants of classical Fermattype equations. In course of time, this line of research has gained momentum, leading to a number of interesting results by many researchers, thereby enriching the field.

Theorem 1.1 ([12]). For any two positive integers m and n with $m \neq n$, the equation

$$f'(z)^n + f(z+c)^m = 1,$$

has no transcendental entire solution with finite order.

Theorem 1.2 ([12]). The finite order transcendental entire solution of

$$f'(z)^2 + f(z+c)^2 = 1,$$

must satisfy $f(z) = \sin(z \pm Bi)$, where B is a constant and $c = 2k\pi$ or $c = (2k+1)\pi$, k is an integer.

²⁰²⁰ Mathematics Subject Classification. 32H30, 35M30.

Key words and phrases. Fermat-type equation; transcendental entire solution; partial differential operator.

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Submitted June 12, 2024. Published October 15, 2024.

In 2017, Gao [5] investigated the existence and form of transcendental entire solutions, for the following systems of Fermat-type equations:

$$f_1'(z)^{n_1} + f_2(z+c)^{m_1} = Q_1(z),$$

$$f_2'(z)^{n_2} + f_1(z+c)^{m_2} = Q_2(z),$$
(1.1)

and

$$f_1'(z)^2 + f_2(z+c)^2 = Q_1(z),$$

$$f_2'(z)^2 + f_1(z+c)^2 = Q_2(z),$$
(1.2)

where $Q_j(z)$, j = 1, 2 are non-zero polynomials in \mathbb{C} . For systems (1.1) and (1.2), the following results were obtained:

Theorem 1.3 ([5]). There does not exist any finite order transcendental entire solutions $(f_1(z), f_2(z))$ of (1.1) if any of the following conditions is satisfied:

(i) $m_1 m_2 > n_1 n_2;$ (ii) $m_j > \frac{n_j}{n_j - 1}, \ j = 1, 2.$

Theorem 1.4 ([5]). Let $(f_1(z), f_2(z))$ be a finite order transcendental entire solution of (1.2) in \mathbb{C} . Then $Q_1(z) = c_{11}c_{12}$, $Q_2(z) = c_{21}c_{22}$ and

$$f_1(z) = \frac{c_{11}e^{az+b_1} - c_{12}e^{-az-b_1}}{2a}, \quad f_2(z) = \frac{c_{21}e^{az+b_2} - c_{21}e^{-az-b_2}}{2a},$$

where $a^4 = 1$, b_1 , b_2 , $c_{kj} (\neq 0)$, k, j = 1, 2 are constants.

In 2018, Xu-Cao [15] investigated the existence and form of transcendental entire solutions of shift-differential equation in \mathbb{C}^2 to obtain the following result.

Theorem 1.5 ([15]). Let $c = (c_1, c_2)$ be a non-zero constant in \mathbb{C}^2 . Then the Fermat-type partial differential equation

$$\left(\frac{\partial f(z_1, z_2)}{\partial z_1}\right)^n + f(z_1 + c_1, z_2 + c_2)^m = 1,$$

does not have a finite order transcendental entire solution whenever m, n are two distinct positive integer.

Theorem 1.6 ([15, 16]). Let $c = (c_1, c_2)$ be a non-zero constant in \mathbb{C}^2 , then each finite order transcendental entire solution of the Fermat-type partial differential equation

$$\left(\frac{\partial f(z_1, z_2)}{\partial z_1}\right)^2 + f(z_1 + c_1, z_2 + c_2)^2 = 1,$$

has the form $f(z_1, z_2) = sin(Az_1 + Bz_2 + H(z_2))$, where A, B are constants in \mathbb{C} satisfying $A^2 = 1$, $e^{i(Ac_1+Bc_2)} = 1$ and $H(z_2)$ is a polynomial in one variable z_2 such that $H(z_2) = H(z_2 + c_2)$. In special case whenever $c_2 \neq 0$, we have $f(z_1, z_2) = sin(Az_1 + Bz_2 + constant)$.

Motivated by the results several authors made contribution to this field; see [1, 14], [17]-[21] and the references therein.

2. FORMULATION OF MAIN PROBLEM AND RELEVANT EXAMPLES

To proceed further, we introduce the following differential-operator

Definition 2.1. The partial differential operator P_{L_n} in \mathbb{C}^n is defined as

$$P_{L_n} = \sum_{|J|=1}^n b_{j_1...j_n} \frac{\partial^{|J|}}{\partial z_1^{j_1} \dots \partial z_n^{j_n}},$$

where $b_{j_1...j_n} \neq 0$ are constants in \mathbb{C} and $J = (j_1, j_2, ..., j_n), |J| = \sum_{t=1}^n j_t$. From now onwards, we use $\vec{z}_n = (z_1, z_2, ..., z_n), \ \vec{c}_n = (c_1, c_2, ..., c_n) \ \text{and} \ \vec{z_n} + \vec{c_n} = (z_1 + c_1, z_2 + c_2, ..., z_n + c_n) \ \text{and} \ \vec{0} = (0, 0, ..., 0)$.

This article is based on exploring existence of finite order transcendental entire solutions in n ($n \ge 1$) dimensional complex field of the equations

$$(P_{L_n}(f_1(\vec{z_n})))^{l_1} + f_2(\vec{z_n} + \vec{c_n})^{k_1} = Q_1(\vec{z_n})$$

$$(P_{L_n}(f_2(\vec{z_n})))^{l_2} + f_1(\vec{z_n} + \vec{c_n})^{k_2} = Q_2(\vec{z_n});$$
(2.1)

where $Q_j(\vec{z_n})$, j=1,2 are two non-zero polynomials in \mathbb{C}^n and are of finite order transcendental entire solution for n=2, i.e. in \mathbb{C}^2 of the equations

$$(P_{L_2}(f_1(\vec{z_2})))^2 + f_2(\vec{z_2} + \vec{c_2})^2 = 1$$

$$(P_{L_2}(f_2(\vec{z_2})))^2 + f_1(\vec{z_2} + \vec{c_2})^2 = 1.$$
(2.2)

Theorem 2.2. Let $\vec{c} = (c_1, c_2, ..., c_n)$ be a non-zero constant in \mathbb{C}^n . Then (2.1) can not have a finite order transcendental entire solution $(f_1(\vec{z_n}), f_2(\vec{z_n}), ..., f_n(\vec{z_n}))$ if the exponents satisfy one of the following two conditions:

- (i) $k_1k_2 > l_1l_2$;
- (ii) $k_t > \frac{l_t}{l_t-1}$ for $l_t \geq 2$, t = 1, 2.

The above theorem motivate us to explore the case $l_t = 1$, and $k_t = 1$; t = 1, 2. In this respect, the following example shows that the solution exists.

Example 2.3. Let $l_1 = 1$, $l_2 = 1$, $k_1 = 1$, $k_2 = 1$, $b_{10} = 1$, $b_{01} = 1$, $b_{11} = 1$, $b_{20} = 1$, $b_{02} = 1$, $Q_1(\vec{z_2}) = 1$, $Q_2(\vec{z_2}) = 1$. Then $f(z) = (f_1(\vec{z_2}), f_2(\vec{z_2}))$, where $f_j(\vec{z_2}) = e^{z_1 + z_2} + 1$, j = 1, 2 is a solution of (2.1) when $e^{c_1 + c_2} = -5$.

For the sake of convenience and to proceed further we us use the following expressions

$$\begin{split} A_1(r,s) &= -b_{10}s + b_{01}r + b_{11}(d_2s - d_1r) + 2b_{20}d_1s - 2b_{02}d_2r, \\ A_2(r,s) &= b_{10}s - b_{01}r + b_{11}(d_2s - d_1r) + 2b_{20}d_1s - 2b_{02}d_2r, \\ B(r,s) &= -b_{11}rs + b_{20}s^2 + b_{02}r^2, \\ D_1(r,s) &= -b_{10}r - b_{01}s + b_{11}rs + b_{20}r^2 + b_{02}s^2, \\ D_2(r,s) &= b_{10}r + b_{01}s + b_{11}rs + b_{20}r^2 + b_{02}s^2, \\ L(r,s) &= d_1r + d_2s, \end{split}$$

where r, s are parameters and d_1 , d_2 are two constants in \mathbb{C} .

Theorem 2.4. Let $(c_1, c_2) \neq (0, 0) \in \mathbb{C}^2$ be a constant and $(f_1(\vec{z_2}), f_2(\vec{z_2}))$ be a finite order transcendental entire solution of (2.2) in \mathbb{C}^2 . Also let $B(c_1, c_2)$,

 $A_1(c_1, c_2)$ and $B(c_1, c_2)$, $A_2(c_1, c_2)$ be nonzero simultaneously. Then $(f_1(\vec{z_2})f_2(\vec{z_2}))$ takes one of the following form:

(A) When

$$D_1(d_1, d_2)D_2(d_1, d_2) = 1, \quad e^{2L(\vec{c_2})} = \frac{D_2(d_1, d_2)}{D_1(d_1, d_2)},$$
$$e^{2(W_1 + W_2)} = -1, \quad e^{W_1 + W_2} = \frac{i}{D_1(d_1, d_2)}e^{-L(\vec{c_2})},$$

we have

$$f_1(\vec{z_2}) = \frac{e^{-L(\vec{z_2}) + L(\vec{c_2}) + W_2} - e^{L(\vec{z_2}) - L(\vec{c_2}) - W_2}}{2i},$$

$$f_2(\vec{z_2}) = \frac{e^{L(\vec{z_2}) - L(\vec{c_2}) + W_1} - e^{-L(\vec{z_2}) + L(\vec{c_2}) - W_1}}{2i},$$

where W_1 , W_2 are two constants in \mathbb{C} .

(B) When

$$\begin{split} D_1(d_1,d_2)D_2(d_1,d_2) &= 1, \quad e^{2L(\vec{c_2})} = -\frac{D_2(d_1,d_2)}{D_1(d_1,d_2)}, e^{2(W_1-W_2)} = 1, \\ e^{W_1-W_2} &= -\frac{i}{D_1(d_1,d_2)}e^{-L(\vec{c_2})}, \\ f_1(\vec{z_2}) &= \frac{e^{L(\vec{z_2})-L(\vec{c_2})+W_2}-e^{-L(\vec{z_2})+L(\vec{c_2})-W_2}}{2i}, \\ f_2(\vec{z_2}) &= \frac{e^{L(\vec{z_2})-L(\vec{c_2})+W_1}-e^{-L(\vec{z_2})+L(\vec{c_2})-W_1}}{2i}, \end{split}$$

where W_1 , W_2 are two constants in \mathbb{C} .

The following examples justify Theorem 2.4.

Example 2.5. Let $b_{10}=-2i,\ b_{01}=i,\ b_{11}=-2i,\ b_{20}=i,\ b_{02}=i,\ d_1=1,\ d_2=1,$ $c_1=\frac{\pi i}{4},\ c_2=\frac{\pi i}{4},\ W_1=\frac{-\pi i}{4},\ W_2=\frac{-\pi i}{4}.$ Then

$$(f_1(\vec{z}), f_2(\vec{z_2})) = \left(\frac{-ie^{-z_1 - z_2 + \frac{\pi i}{4}} + ie^{z_1 + z_2 - \frac{\pi i}{4}}}{2}, -\frac{e^{-z_1 - z_2 + \frac{\pi i}{4}} + e^{z_1 + z_2 - \frac{\pi i}{4}}}{2}\right)$$

is a solution of (2.2).

Example 2.6. Let $b_{10}=2i,\ b_{01}=i,\ b_{11}=2i,\ b_{20}=i,\ b_{02}=i,\ d_1=1,\ d_2=-1,\ c_1=\frac{\pi i}{4},\ c_2=-\frac{\pi i}{4},\ W_1=\frac{\pi i}{4},\ W_2=\frac{\pi i}{4}.$ Then

$$(f_1(\vec{z}), f_2(\vec{z_2})) = \left(\frac{ie^{-z_1 + z_2 - \frac{\pi i}{4}} - ie^{z_1 - z_2 + \frac{\pi i}{4}}}{2}, -\frac{e^{-z_1 + z_2 - \frac{\pi i}{4}} + e^{z_1 - z_2 + \frac{\pi i}{4}}}{2}\right)$$

is a solution of (2.2).

Example 2.7. Let $b_{10} = 2i$, $b_{01} = i$, $b_{11} = 2i$, $b_{20} = i$, $b_{02} = i$, $d_1 = 1$, $d_2 = -1$, $c_1 = 2\pi i$, $c_2 = -2\pi i$, $W_1 = 0$, $W_2 = 0$. Then

$$(f_1(\vec{z_2}), f_2(\vec{z_2})) = \left(\frac{ie^{-z_1+z_2} - ie^{z_1-z_2}}{2}, \frac{ie^{-z_1+z_2} - ie^{z_1-z_2}}{2}\right)$$

is a solution of (2.2).

Example 2.8. Let $b_{10} = 2i$, $b_{01} = i$, $b_{11} = 2i$, $b_{20} = i$, $b_{02} = i$, $d_1 = 1$, $d_2 = -1$, $c_1 = 1$, $c_2 = 1$, $W_1 = 0$, $W_2 = 0$. Then

$$(f_1(\vec{z}), f_2(\vec{z_2})) = \left(\frac{ie^{-z_1+z_2} - ie^{z_1-z_2}}{2}, \frac{ie^{-z_1+z_2} - ie^{z_1-z_2}}{2}\right)$$

is a solution of (2.2).

Example 2.9. Let $b_{10} = i$, $b_{01} = 2i$, $b_{11} = 2i$, $b_{20} = i$, $b_{02} = i$, $d_1 = 1$, $d_2 = -1$, $c_1 = \pi i$, $c_2 = -\pi i$, $W_1 = \frac{\pi i}{2}$, $W_2 = -\frac{\pi i}{2}$. Then

$$(f_1(\vec{z_2}), f_2(\vec{z_2})) = \left(\frac{-ie^{-z_1+z_2} + ie^{z_1-z_2}}{2}, \frac{ie^{-2z_1+z_2} - ie^{z_1-z_2}}{2}\right)$$

is a solution of (2.2).

Example 2.10. Let $b_{10} = -2i$, $b_{01} = -i$, $b_{11} = 2i$, $b_{20} = i$, $b_{02} = i$, $d_1 = 1$, $d_2 = -1$, $c_1 = 1$, $c_2 = 1$, $W_1 = 2\pi i$, $W_2 = \pi i$. Then

$$(f_1(\vec{z}), f_2(\vec{z_2})) = \left(\frac{-ie^{-z_1+z_2} + ie^{z_1-z_2}}{2}, \frac{ie^{-z_1+z_2} - ie^{z_1-z_2}}{2}\right)$$

is a solution of (2.2).

Example 2.11. Let $b_{10} = 1$, $b_{01} = 1$, $b_{11} = 1$, $b_{20} = 1$, $b_{02} = 1$, $d_1 = 1$, $d_2 = -1$, $c_1 = c_2 = 1$, $W_1 = \frac{\pi i}{4}$, $W_2 = \frac{\pi i}{4}$. Then

$$(f_1(\vec{z}), f_2(\vec{z_2})) = \Big(\frac{e^{-z_1 + z_2 - \frac{\pi i}{4}} + e^{z_1 - z_2 + \frac{\pi i}{4}}}{2}, \frac{ie^{-z_1 + z_2 - \frac{\pi i}{4}} - ie^{z_1 - z_2 + \frac{\pi i}{4}}}{2}\Big)$$

is a solution of (2.2).

Example 2.12. Let $b_{10} = 1$, $b_{01} = 2$, $b_{11} = 2$, $b_{20} = 1$, $b_{02} = 1$, $d_1 = 2$, $d_2 = -1$, $c_1 = 2\pi i$, $c_2 = 2\pi i$, $W_1 = \frac{\pi i}{4}$, $W_2 = \frac{\pi i}{4}$. Then

$$(f_1(\vec{z}),f_2(\vec{z_2})) = \left(\frac{e^{-2z_1+z_2-\frac{\pi i}{4}}+e^{2z_1-z_2+\frac{\pi i}{4}}}{2},\frac{ie^{-2z_1+z_2-\frac{\pi i}{4}}-ie^{2z_1-z_2+\frac{\pi i}{4}}}{2}\right)$$

is a solution of (2.2).

Example 2.13. Let $b_{10} = 1$, $b_{01} = 2$, $b_{11} = 2$, $b_{20} = 1$, $b_{02} = 1$, $d_1 = 2i$, $d_2 = -i$, $c_1 = i$, $c_2 = 2i$, $W_1 = -\frac{\pi i}{4}$, $W_2 = -\frac{\pi i}{4}$. Then

$$(f_1(\vec{z}),f_2(\vec{z_2})) = \Big(-\frac{e^{-2iz_1+iz_2+\frac{\pi i}{4}}+e^{2iz_1-iz_2-\frac{\pi i}{4}}}{2},\frac{ie^{-2iz_1+iz_2+\frac{\pi i}{4}}-ie^{2iz_1-iz_2-\frac{\pi i}{4}}}{2} \Big)$$

is a solution of (2.2).

Example 2.14. Let $b_{10} = 1$, $b_{01} = 2$, $b_{11} = 2$, $b_{20} = 1$, $b_{02} = 1$, $d_1 = 2i$, $d_2 = -i$, $c_1 = 2\pi$, $c_2 = 2\pi$, $W_1 = -\frac{\pi i}{4}$, $W_2 = -\frac{\pi i}{4}$, Then

$$(f_1(\vec{z}),f_2(\vec{z_2})) = \Big(-\frac{e^{-2iz_1+iz_2+\frac{\pi i}{4}}+e^{2iz_1-iz_2-\frac{\pi i}{4}}}{2},\frac{ie^{-2iz_1+iz_2+\frac{\pi i}{4}}-ie^{2iz_1-iz_2-\frac{\pi i}{4}}}{2} \Big)$$

is a solution of (2.2).

Example 2.15. Let $b_{10} = 1$, $b_{01} = 2$, $b_{11} = 2$, $b_{20} = 1$, $b_{02} = 1$, $d_1 = 2$, $d_2 = -1$, $c_1 = \frac{\pi i}{2}$, $c_2 = -\frac{\pi i}{2}$, $W_1 = \frac{\pi i}{2}$, $W_2 = \frac{\pi i}{2}$. Then

$$(f_1(\vec{z}), f_2(\vec{z_2})) = \left(\frac{e^{-2z_1+z_2} + e^{2z_1-z_2}}{2}, \frac{e^{-2z_1+z_2} + e^{2z_1-z_2}}{2}\right)$$

is a solution of (2.2).

Example 2.16. Let $b_{10} = b_{01} = -\frac{6}{5}$, $b_{11} = \frac{1}{5}$, $b_{20} = \frac{8}{5}$, $b_{02} = \frac{4}{5}$, $d_1 = 1$, $d_2 = 1$, $c_1 = 1$, $c_2 = \log(\frac{1}{5}) - 1$, $W_1 = W_2 = \frac{\pi i}{4}$. Then

$$(f_1(\vec{z}), f_2(\vec{z_2})) = \left(\frac{\frac{1}{5}e^{-z_1 - z_2 + \frac{\pi i}{4}} - 5e^{z_1 + z_2 - \frac{\pi i}{4}}}{2i}, \frac{5e^{z_1 + z_2 + \frac{\pi i}{4}} - \frac{1}{5}e^{-z_1 - z_2 - \frac{\pi i}{4}}}{2i}\right)$$

is a solution of (2.2).

Corollary 2.17. Let $(f_1(\vec{z_2}), f_2(\vec{z_2}))$ be a transcendental entire function of order properly greater than one with $B(c_1, c_2)$, $A_1(c_1, c_2)$ and $A_2(c_1, c_2)$ are not zero simultaneously. Then $(f_1(\vec{z_2}), f_2(\vec{z_2}))$ can not be a solution of (2.2).

3. Lemmas

We assume that the readers are familiar with the basic notations of the Nevanlinna theory such as N(r,f), $N(r,\frac{1}{f})$, m(r,f), T(r,f) in complex variable [7]. For several complex variables we refer to [10] and the references therein. By S(r,f) we will mean any quantity satisfying S(r,f) = o(T(r,f)), $r \to \infty$, outside possibly an exceptional set of finite logarithmic measure. Based on the notations, the following lemmas will play important role in proving our theorems.

Lemma 3.1 ([13]). For each entire function F in \mathbb{C}^n , $F(\vec{0}) \neq \vec{0}$ and put $\rho(n_F) = \rho < \infty$. Then there exists a canonical function f_F and a function $g_F \in \mathbb{C}^n$ such that $F(z) = f_{F(z)}e^{g_F(\vec{z}_n)}$. For special case n = 1, f_F is the canonical product of Weierstrass. Here $\rho(n_f)$ denotes the order of the counting function of zeros of F.

Lemma 3.2 ([22, 2]). Let f(z) be a non-constant meromorphic function in \mathbb{C}^n and let $I = (i_1, \ldots, i_n)$ be a multi index with length $|I| = \sum_{j=1}^n i_j$. Assume that $T(r_0, f) \geq e$ for some r_0 . Then

$$m\left(r, \frac{\partial^I f}{f}\right) = S(r, f),$$

holds for all $r \geq r_0$, outside a set $E \subset (0, +\infty)$ of finite logarithmic measure $\int_E \frac{dt}{t} < \infty$, where $\partial^I f = \frac{\partial^I f}{\partial z_1^{i_1} ... \partial z_n^{i_n}}$.

Lemma 3.3 ([9]). Let $f_j(\not\equiv 0)$, j=1,2,3 be meromorphic function in \mathbb{C}^n such that f_1 is not constant, $f_1+f_2+f_3=1$ and

$$\sum_{i=1}^{3} \left\{ N_2 \left(r, \frac{1}{f_j} \right) + 2\overline{N}(r, f_j) \right\} < \lambda T(r, f_1) + o(\log^+ T(r, f_1)),$$

for all r outside possibly a set with finite logarithmic measure, where $\lambda < 1$ is a positive number, then either $f_2 \equiv 1$ or $f_3 \equiv 1$.

Lemma 3.4 ([3]). Let $f(\vec{z}_n)$ be a non-constant meromorphic function with finite order in \mathbb{C}^n such that $f(\vec{0}) \neq 0, \infty$ and let $\epsilon > 0$. Then, for $\vec{c}_n \in \mathbb{C}^n$,

$$m\left(r, \frac{f(\vec{z}_n)}{f(\vec{z}_n + \vec{c}_n)}\right) + m\left(r, \frac{f(\vec{z}_n + \vec{c}_n)}{f(\vec{z}_n)}\right) = S(r, f),$$

holds for all $r \geq r_0$, outside a set $E \subset (0, +\infty)$ of finite logarithmic measure $\int_E \frac{dt}{t} < \infty$.

Lemma 3.5 ([9, Lemma 3.1]). Suppose that $a_0(\vec{z}_m), a_1(\vec{z}_m), \ldots, a_n(\vec{z}_m), n \geq 1$, are meromorphic in \mathbb{C}^m and $g_0(\vec{z}_m), g_2(\vec{z}_m), \ldots, g_n(\vec{z}_m)$ are entire in \mathbb{C}^m . $g_j(\vec{z}_m) - g_k(\vec{z}_m)$ are non-constant for $0 \leq j < k \leq n$. If

$$\sum_{j=0}^{n} a_j(\vec{z}_m) e^{g_j(\vec{z}_m)} \equiv 0$$

and $T(r, a_j) = o(T(r)), j = 0, 1, 2 \dots, n,$

$$T(r) = \min_{0 \le j < k \le n} T(r, e^{g_k - g_j}),$$

then $a_i \equiv 0$.

4. Proof of the main results

Proof of Theorem 2.2. Let $(f_1(\vec{z_n}), f_2(\vec{z_n}), \dots, f_n(\vec{z_n}))$ be a finite order transcendental entire solution of (2.1) in \mathbb{C}^n . We consider the following 2 cases:

Case 1: Let $k_1k_2 > l_1l_2$. Using Lemma 3.4, we have that

$$m\left(r, \frac{f_j(\vec{z_n})}{f_j(\vec{z_n} + \vec{c_n})}\right) = S(r, f_j), \tag{4.1}$$

holds for all r > 0 outside a possible set $E_j \subset [1, +\infty)$, j = 1, 2, ..., n of finite logarithmic measure $\int_{E_j} \frac{dt}{t} < \infty$. Clearly, we have the following

$$T(r, f_{j}(\vec{z_{n}})) = m(r, f_{j}(\vec{z_{n}})),$$

$$\leq m\left(r, \frac{f_{j}(\vec{z_{n}})}{f_{j}(\vec{z_{n}} + \vec{c_{n}})} \times f_{j}(\vec{z_{n}} + \vec{c_{n}})\right),$$

$$\leq m\left(r, \frac{f_{j}(\vec{z_{n}})}{f_{j}(\vec{z_{n}} + \vec{c_{n}})}\right) + m\left(r, f_{j}(\vec{z_{n}} + \vec{c_{n}})\right) + \log 2,$$

$$= m(r, f_{j}(\vec{z_{n}} + \vec{c_{n}})) + \log 2 + S(r, f_{j}),$$

$$= T(r, f_{j}(\vec{z_{n}} + \vec{c_{n}})) + \log 2 + S(r, f_{j}),$$

$$= T(r, f_{j}(\vec{z_{n}} + \vec{c_{n}})) + \log 2 + S(r, f_{j}),$$

$$(4.2)$$

for all $r \notin E_1 \cup E_2$. Applying Valliron Mohon'ko theorem in several complex variables [8] we have

$$k_{1}T(r, f_{2}(\vec{z_{n}})) \leq k_{1}T(r, f_{2}(\vec{z_{n}} + \vec{c_{n}})) + S(r, f_{2})$$

$$\leq T(r, f_{2}(\vec{z_{n}} + \vec{c_{n}}))^{k_{1}}) + S(r, f_{2}),$$

$$= T\left(r, (P_{L_{n}}(f_{1}(\vec{z_{n}})))^{l_{1}} - Q_{1}(\vec{z_{n}})\right) + S(r, f_{2}),$$

$$= l_{1}T\left(r, P_{L_{n}}(f_{1}(\vec{z_{n}}))\right) + S(r, f_{1}) + S(r, f_{2}),$$

$$= l_{1}m\left(r, P_{L_{n}}(f_{1}(\vec{z_{n}}))\right) + S(r, f_{1}) + S(r, f_{2}),$$

$$\leq l_{1}\left[m\left(r, \frac{P_{L_{n}}(f_{1}(\vec{z_{n}}))}{f_{1}(\vec{z_{n}})}\right) + m(r, f_{1}(\vec{z_{n}})) + \log 2\right]$$

$$+ S(r, f_{1}) + S(r, f_{2}),$$

$$= l_{1}T(r, f_{1}(\vec{z_{n}}) + S(r, f_{1}) + S(r, f_{2}),$$

$$(4.3)$$

i.e. from (4.3) wee obtain

$$(k_1 + o(1))T(r, f_2(\vec{z_n})) \le (l_1 + o(1))T(r, f_1(\vec{z_n})), \quad r \notin E_1.$$
(4.4)

Similarly, we obtain

$$(k_2 + o(1))T(r, f_2(\vec{z_n})) \le (l_2 + o(1))T(r, f_1(\vec{z_n})), \quad r \notin E_2.$$
 (4.5)

From (4.4) and (4.5) clearly we have a contradiction.

Case 2: Let $k_t > \frac{l_t}{l_t-1}$, $l_t \ge 2$, t = 1, 2. Using the Nevanlinna second main theorem, from (2.1) we obtain

$$\begin{split} &(l_{1}-1)T\left(r,P_{L_{n}}(f_{1}(\vec{z_{n}}))\right) \\ &\leq \overline{N}\left(r,P_{L_{n}}(f_{1}(\vec{z_{n}}))\right) + \overline{N}\left(r,\frac{1}{\left(P_{L_{n}}(f_{1}(\vec{z_{n}}))\right)^{l_{1}} - Q_{1}(\vec{z_{n}})}\right) + S\left(r,P_{L_{n}}(f_{1})\right), \\ &\leq \overline{N}\left(r,\frac{1}{f_{2}(\vec{z_{n}}+\vec{c_{n}})}\right) + S(r,f_{1}), \\ &\leq T(r,f_{2}(\vec{z_{n}}+\vec{c_{n}})) + S(r,f_{1}), \\ &\leq T(r,f_{2}(\vec{z_{n}})) + S(r,f_{1}) + S(r,f_{2}). \end{split} \tag{4.6}$$

Proceeding with the similar arguments, from the second equation we obtain

$$(l_2 - 1)T(r, P_{L_n} f_2((\vec{z_n})) \le T(r, f_1(\vec{z_n})) + S(r, f_1) + S(r, f_2). \tag{4.7}$$

From the first equation of (2.1) and using Valliron Mohon'ko theorem in several complex variables [8] we obtain

$$k_1 T(r, f_2(\vec{z_n} + \vec{c_n})) = T(r, (P_{L_n}(f_1(\vec{z_n})))^{l_1} - Q_1(\vec{z_n})) + S(r, f_1)$$

$$\leq l_1 T(r, P_{L_n}(f_1(\vec{z_n}))) + S(r, f_1) + S(r, f_1).$$
(4.8)

Proceeding, in the similar way from the second equation of (2.1) we obtain

$$k_2T(r, f_1(\vec{z_n} + \vec{c_n})) \le l_2T(r, P_{L_n}(f_2(\vec{z_n})) + S(r, f_1) + S(r, f_2).$$
 (4.9)

From (4.6)-(4.9) we obtain

$$\left(k_1 - \frac{l_1}{l_1 - 1} + o(1)\right) T(r, f_2(\vec{z_n})) \le S(r, f_1),$$

$$\left(k_2 - \frac{l_2}{l_2 - 1} + o(1)\right) T(r, f_1(\vec{z_n})) \le S(r, f_2).$$

Since $(f_1(\vec{z_n}), f_2(\vec{z_n}), \dots, f_n(\vec{z_n}))$ is a transcendental entire function, we obtain

$$\left(k_1 - \frac{l_1}{l_1 - 1} + o(1)\right) \left(k_2 - \frac{l_2}{l_2 - 1} + o(1)\right) \le 0.$$

Since $k_t > \frac{l_t}{l_t-1}$, t=1,2, we have a contradiction. The proof of Theorem 2.2 is complete

The following expression is used several times to prove the next theorem.

$$\begin{split} M_{m,u}(p) &= b_{10} \frac{\partial p_m(\vec{z_2})}{\partial z_1} + b_{01} \frac{\partial p_m(\vec{z_2})}{\partial z_2} + b_{11} \Big\{ \frac{\partial^2 p_m(\vec{z_2})}{\partial z_1 \partial z_2} \\ &+ (-1)^{u-1} \frac{\partial p_m(\vec{z_2})}{\partial z_1} \frac{\partial p_m(\vec{z_2})}{\partial z_2} \Big\} + b_{20} \Big\{ \frac{\partial^2 p_m(\vec{z_2})}{\partial z_1^2} \\ &+ (-1)^{u-1} \Big(\frac{\partial p_m(\vec{z_2})}{\partial z_1} \Big)^2 \Big\} + b_{02} \Big\{ \frac{\partial^2 p_m(\vec{z_2})}{\partial z_2^2} + (-1)^{u-1} \Big(\frac{\partial p_m(\vec{z_2})}{\partial z_2} \Big)^2 \Big\}, \end{split}$$

for m, u = 1, 2.

Proof of Theorem 2.4. Let $(f_1(\vec{z_2}), f_2(\vec{z_2}))$ be a pair of finite order transcendental entire solution of (2.2) in \mathbb{C}^2 . Clearly, system (2.2) can be re-written as follows

$$\begin{aligned}
&\{P_{L_2}(f_1(\vec{z_2})) + if_2(\vec{z_2} + \vec{c_2})\}\{P_{L_2}(f_1(\vec{z_2})) - if_2(\vec{z_2} + \vec{c_2})\} = 1, \\
&\{P_{L_2}(f_2(\vec{z_2}) + if_1(\vec{z_2} + \vec{c_2})\}\{P_{L_2}(f_2(\vec{z_2})) - if_1(\vec{z_2} + \vec{c_2})\} = 1.
\end{aligned} (4.10)$$

Now using Lemma 3.1, from (4.10) we obtain

$$P_{L_{2}}(f_{1}(\vec{z_{2}})) + if_{2}(\vec{z_{2}} + \vec{c_{2}}) = e^{p_{1}(\vec{z_{2}})},$$

$$P_{L_{2}}(f_{1}(\vec{z_{2}})) - if_{2}(\vec{z_{2}} + \vec{c_{2}}) = e^{-p_{1}(\vec{z_{2}})},$$

$$P_{L_{2}}(f_{2}(\vec{z_{2}})) + if_{1}(\vec{z_{2}} + \vec{c_{2}}) = e^{p_{2}(\vec{z_{2}})},$$

$$P_{L_{2}}(f_{2}(\vec{z_{2}}) - if_{1}(\vec{z_{2}} + \vec{c_{2}}) = e^{-p_{2}(\vec{z_{2}})},$$

$$(4.11)$$

where $p_1(\vec{z_2})$, $p_2(\vec{z_2})$ are two non-constant polynomials in \mathbb{C}^2 . By an easy computation from (4.11), we obtain

$$P_{L_{2}}(f_{1}(\vec{z_{2}})) = \frac{e^{p_{1}(\vec{z_{2}})} + e^{-p_{1}(\vec{z_{2}})}}{2},$$

$$f_{2}(\vec{z_{2}} + \vec{c_{2}}) = \frac{e^{p_{1}(\vec{z_{2}})} - e^{-p_{1}(\vec{z_{2}})}}{2i},$$

$$P_{L_{2}}(f_{2}(\vec{z_{2}})) = \frac{e^{p_{2}(\vec{z_{2}})} + e^{-p_{2}(\vec{z_{2}})}}{2},$$

$$f_{1}(\vec{z_{2}} + \vec{c_{2}}) = \frac{e^{p_{2}(\vec{z_{2}})} - e^{-p_{2}(\vec{z_{2}})}}{2i}.$$

$$(4.12)$$

Combining the first and the last equations, and the second and the third equations of (4.12) we obtain respectively

$$-iM_{2.1}e^{p_1(\vec{z_2}+\vec{c_2})+p_2(\vec{z_2})} - iM_{2.2}e^{p_1(\vec{z_2}+\vec{c_2})-p_2(\vec{z_2})} - e^{2p_1(\vec{z_2}+\vec{c_2})} = 1, \tag{4.13}$$

and

$$-iM_{1,1}e^{p_2(\vec{z_2}+\vec{c_2})+p_1(\vec{z_2})}-iM_{1,2}e^{p_2(\vec{z_2}+\vec{c_2})-p_1(\vec{z_2})}-e^{2p_2(\vec{z_2}+\vec{c_2})}=1. \hspace{1.5cm} (4.14)$$

Now taking into consideration equation (4.13), we discuss the following possibilities:

- (i) Let $M_{2,1} \equiv 0$, $M_{2,2} \equiv 0$. Then we have $-e^{2p_1(\vec{z_2}+\vec{c_2})} = 1$, which shows that $p_1(\vec{z_2})$ is a constant polynomial, a contradiction.
- (ii) Let $M_{2,1} \equiv 0$ and $M_{2,2} \not\equiv 0$. Then we have

$$-iM_{2,2}e^{p_1(\vec{z_2}+\vec{c_2})-p_2(\vec{z_2})} - e^{2p_1(\vec{z_2}+\vec{c_2})} = 1. (4.15)$$

Since $p_1(\vec{z_2})$ is a non-constant polynomial, (4.15) implies that $p_1(\vec{z_2} + \vec{c_2}) - p_2(\vec{z_2})$ is also non-constant. We claim that $-p_2(\vec{z_2}) - p_1(\vec{z_2} + \vec{c_2})$ is also non-constant. On the contrary, let $-p_2(\vec{z_2}) - p_1(\vec{z_2} + \vec{c_2}) = A_1'$, where A_1' is a constant in \mathbb{C} .

Then from (4.15) we obtain

$$-iM_{2,2}e^{A'_1+2p_1(\vec{z_2}+\vec{c_2})} - e^{2p_1(\vec{z_2}+\vec{c_2})} = 1,$$

i.e. $(iM_{2,2}e^{A'_1}+1)e^{2p_1(\vec{z_2}+\vec{c_2})} = -1.$

Then we have $p_1(\vec{z_2})$ is a constant polynomial, a contradiction. Clearly, we can rewrite (4.15) as

$$-iM_{2,2}e^{-p_2(\vec{z_2})} - e^{p_1(\vec{z_2} + \vec{c_2})} - e^{-p_1(\vec{z_2} + \vec{c_2})} = 0.$$
(4.16)

Now applying Lemma 3.5 in (4.16) we have $M_{2,2} \equiv 0$, a contradiction.

(iii) Let $M_{2,1} \not\equiv 0$ and $M_{2,2} \equiv 0$. Then proceeding in the similar way as done in case (ii) we obtain a contradiction.

So we must have $M_{2,1} \not\equiv 0$ and $M_{2,2} \not\equiv 0$. Using similar arguments from (4.14) we obtain $M_{1,1} \not\equiv 0$ and $M_{1,2} \not\equiv 0$. Hence using Lemma 3.3, in (4.13) and (4.14) we obtain

$$-iM_{2,1}e^{p_1(\vec{z_2}+\vec{c_2})+p_2(\vec{z_2})} \equiv 1 \text{ or } -iM_{2,2}e^{p_1(\vec{z_2}+\vec{c_2})-p_2(\vec{z_2})} \equiv 1;$$

$$-iM_{1,1}e^{p_2(\vec{z_2}+\vec{c_2})+p_1(\vec{z_2})} \equiv 1 \text{ or } -iM_{1,2}e^{p_2(\vec{z_2}+\vec{c_2})-p_1(\vec{z_2})} \equiv 1,$$

respectively.

Now we consider the following four cases:

Case 1:

$$-iM_{2,1}e^{p_1(\vec{z_2}+\vec{c_2})+p_2(\vec{z_2})} \equiv 1,$$

$$-iM_{1,1}e^{p_2(\vec{z_2}+\vec{c_2})+p_1(\vec{z_2})} \equiv 1.$$

Clearly we have $p_1(\vec{z_2} + \vec{c_2}) + p_2(\vec{z_2}) \equiv \eta_1$, $p_2(\vec{z_2} + \vec{c_2}) + p_1(z) \equiv \eta_2$, where η_1 , η_2 are two constants in \mathbb{C} . Then we have $p_1(\vec{z_2}) = L(\vec{z_2}) + H(s) + W_1$, $p_2(\vec{z_2}) = -L(\vec{z_2}) - H(s) + W_2$, where W_1 , W_2 are two constants in \mathbb{C} , H(s) is a polynomial in $s = c_2 z_1 - c_1 z_2$. Now combining with (4.13) and (4.14) we obtain

$$\begin{aligned} b_{10}\left(-d_{1}-H'(s)c_{2}\right)+b_{01}\left(-d_{2}+H'(s)c_{1}\right)\\ +b_{11}\left\{H''(s)c_{1}c_{2}+\left(-d_{1}-H'(s)c_{2}\right)\left(-d_{2}+H'(s)c_{1}\right)\right\}\\ +b_{20}\left\{-H''(s)c_{2}^{2}+\left(-d_{1}-H'(s)c_{2}\right)^{2}\right\}\\ +b_{02}\left\{-H''(s)c_{1}^{2}+\left(-d_{2}+H'(s)c_{1}\right)^{2}\right\}e^{L(\vec{c_{2}})+W_{1}+W_{2}}\equiv i,\\ b_{10}\left(d_{1}+H'(s)c_{2}\right)+b_{01}\left(d_{2}-H'(s)c_{1}\right)\\ +b_{11}\left\{-H''(s)c_{1}c_{2}+\left(d_{1}+H'(s)c_{2}\right)\left(d_{2}-H'(s)c_{1}\right)\right\}+b_{20}\left\{H''(s)c_{2}^{2}\right.\\ +\left(d_{1}+H'(s)c_{2}\right)^{2}\right\}+b_{02}\left\{H''(s)c_{1}^{2}+\left(d_{2}-H'(s)c_{1}\right)^{2}\right\}e^{-L(\vec{c_{2}})+W_{1}+W_{2}}\equiv i,\\ b_{10}\left(-d_{1}-H'(s)c_{2}\right)+b_{01}\left(-d_{2}+H'(s)c_{1}\right)\\ +b_{11}\left\{H''(s)c_{1}c_{2}-\left(-d_{1}-H'(s)c_{2}\right)\left(-d_{2}+H'(s)c_{1}\right)\right\}\\ +b_{20}\left\{-H''(s)c_{1}^{2}-\left(-d_{2}+H'(s)c_{1}\right)^{2}\right\}e^{-L(\vec{c_{2}})-W_{1}-W_{2}}\equiv i,\\ b_{10}\left(d_{1}+H'(s)c_{2}\right)+b_{01}\left(d_{2}-H'(s)c_{1}\right)\\ +b_{11}\left\{-H''(s)c_{1}c_{2}-\left(d_{1}+H'(s)c_{2}\right)\left(d_{2}-H'(s)c_{1}\right)\right\}\\ +b_{20}\left\{H''(s)c_{2}^{2}-\left(d_{1}+H'(s)c_{2}\right)^{2}\right\}\\ +b_{02}\left\{H'''(s)c_{1}^{2}-\left(d_{2}-H'(s)c_{1}\right)^{2}\right\}e^{L(\vec{c_{2}})-W_{1}-W_{2}}\equiv i.\end{aligned}$$

We note that coefficient of H'(s) of the first, second, third, and fourth equations are $A_1(c_1,c_2)$, $A_2(c_1,c_2)$, $-A_2(c_1,c_2)$, and $-A_1(c_1,c_2)$ respectively. Also, coefficients of $H'(s)^2$ of the first, second, third, and fourth equations are $B(c_1,c_2)$, $B(c_1,c_2)$, $-B(c_1,c_2)$, and $-B(c_1,c_2)$ respectively. Further, the coefficients of H''(s) of the first, second, third, and fourth equations are $-B(c_1,c_2)$, $B(c_1,c_2)$, $-B(c_1,c_2)$, and $B(c_1,c_2)$ respectively.

Then (4.17) reduces to

$$[D_{1}(d_{1}, d_{2}) + A_{1}(c_{1}, c_{2})H'(s) + B(c_{1}, c_{2})\{H'(s)^{2} - H''(s)\}]$$

$$\times e^{L(\vec{c_{2}}) + W_{1} + W_{2}} \equiv i,$$

$$[D_{2}(d_{1}, d_{2}) + A_{2}(c_{1}, c_{2})H'(s) + B(c_{1}, c_{2})\{H'(s)^{2} + H''(s)\}]$$

$$\times e^{-L(\vec{c_{2}}) + W_{1} + W_{2}} \equiv i,$$

$$[-D_{2}(d_{1}, d_{2}) - A_{2}(c_{1}, c_{2})H'(s) - B(c_{1}, c_{2})\{H'(s)^{2} + H''(s)\}]$$

$$\times e^{-L(\vec{c_{2}}) - W_{1} - W_{2}} \equiv i,$$

$$[-D_{1}(d_{1}, d_{2}) - A_{1}(c_{1}, c_{2})H'(s) - B(c_{1}, c_{2})\{H'(s)^{2} - H''(s)\}]$$

$$\times e^{L(\vec{c_{2}}) - W_{1} - W_{2}} = i$$

$$(4.18)$$

Taking into consideration the first and fourth equations of (4.18), we have the following:

- (a) $A_1(c_1, c_2) \neq 0$, $B(c_1, c_2) \neq 0$, then degree of $H(s) \leq 1$.
- (b) $A_1(c_1, c_2) = 0$, $B(c_1, c_2) \neq 0$, then degree of $H(s) \leq 1$.
- (c) $A_1(c_1, c_2) \neq 0$, $B(c_1, c_2) = 0$, then degree of $H(s) \leq 1$.
- (d) $A_1(c_1, c_2) = 0$, $B(c_1, c_2) = 0$, then degree of H(s) can be any finite number.

Now using the first assumption of Theorem 2.4, i.e. $B(c_1, c_2)$ and $A_1(c_1, c_2)$ are not zero simultaneously, we obtain degree of $H(s) \leq 1$. Since under $H(s) \leq 1$; $p_1(\vec{z_2})$, $p_2(\vec{z_2})$ both become linear polynomials, without loss of generality we can consider $H(s) \equiv 0$. Then from first and fourth equations of (4.18) we must have

$$D_1(d_1, d_2)e^{L(\vec{c_2}) + W_1 + W_2} = i,$$

$$-D_1(d_1, d_2)e^{L(\vec{c_2}) - W_1 - W_2} = i.$$
(4.19)

Let us consider the second and third equations of (4.18). We have the following 4 possibilities:

- (e) $A_2(c_1, c_2) \neq 0$, $B(c_1, c_2) \neq 0$, degree of $H(s) \leq 1$.
- (f) $A_2(c_1, c_2) = 0$, $B(c_1, c_2) \neq 0$, degree of $H(s) \leq 1$.
- (g) $A_2(c_1, c_2) \neq 0$, $B(c_1, c_2) = 0$, degree of $H(s) \leq 1$.
- (h) $A_2(c_1, c_2) = 0$, $B(c_1, c_2) = 0$, degree of H(s) is arbitrary finite number.

Using the assumption of Theorem 2.4, which is $B(c_1, c_2)$ and $A_2(c_1, c_2)$ are not zero simultaneously, we must have $\deg(H(s) \leq 1$. Since $p_1(\vec{z_2})$, $p_2(\vec{z_2})$ becomes a linear polynomial, without any loss of generality we consider $H(s) \equiv 0$. Then from second and third equations of (4.18) we obtain

$$D_2(d_1, d_2)e^{-L(\vec{c_2}) + W_1 + W_2} = i,$$

$$-D_2(d_1, d_2)e^{-L(\vec{c_2}) - W_1 - W_2} = i.$$
(4.20)

Considering all conditions such that degree of $H(s) \leq 1$ i.e. $B(c_1, c_2)$, $A_1(c_1, c_2)$ and $B(c_1, c_2)$, $A_2(c_1, c_2)$ are not zero simultaneously, from (4.19) and (4.20) we have

$$D_1(d_1, d_2)D_2(d_1, d_2) = 1, \quad e^{2L(\vec{c_2})} = \frac{D_2(d_1, d_2)}{D_1(d_1, d_2)},$$
$$e^{2(W_1 + W_2)} = -1, \quad e^{W_1 + W_2} = \frac{i}{D_1(d_1, d_2)}e^{-L(\vec{c_2})}.$$

The form of the solution is

$$f_1(\vec{z_2}) = \frac{e^{-L(\vec{z_2}) + L(\vec{c_2}) + W_2} - e^{L(\vec{z_2}) - L(\vec{c_2}) - W_2}}{2i},$$

$$f_2(\vec{z_2}) = \frac{e^{L(\vec{z_2}) - L(\vec{c_2}) + W_1} - e^{-L(\vec{z_2}) + L(\vec{c_2}) - W_1}}{2i}.$$

Case 2: Let

$$-iM_{2,1}e^{p_1(\vec{z_2}+\vec{c_2})+p_2(\vec{z_2})} \equiv 1,$$

$$-iM_{1,2}e^{p_2(\vec{z_2}+\vec{c_2})-p_1(\vec{z_2})} \equiv 1.$$

Clearly we have $p_1(\vec{z_2} + \vec{c_2}) + p_2(\vec{z_2}) \equiv \eta_1$, $p_2(\vec{z_2} + \vec{c_2}) - p_1(\vec{z_2}) \equiv \eta_2$, where η_1 , η_2 are two constants in \mathbb{C} . Then by easy computation we obtain $p_1(\vec{z_2} + 2\vec{c_2}) + p_1(\vec{z_2}) \equiv \eta_1 - \eta_2$, which contradicts that $p_1(\vec{z_2})$ is a non-constant polynomial.

Case 3: Let

$$-iM_{2,2}e^{p_1(\vec{z_2}+\vec{c_2})-p_2(\vec{z_2})} \equiv 1,$$

$$-iM_{1,1}e^{p_2(\vec{z_2}+\vec{c_2})+p_1(\vec{z_2})} \equiv 1.$$

Then by using similar arguments as in Case 2, we obtain a contradiction.

Case 4: Let

$$-iM_{2,2}e^{p_1(\vec{z_2}+\vec{c_2})-p_2(\vec{z_2})} \equiv 1,$$

$$-iM_{1,2}e^{p_2(\vec{z_2}+\vec{c_2})-p_1(\vec{z_2})} \equiv 1.$$

Then clearly we have $p_1(\vec{z_2} + \vec{c_2}) - p_2(\vec{z_2}) \equiv \eta_1$, $p_2(\vec{z_2} + \vec{c_2}) - p_1(\vec{z_2}) \equiv \eta_2$, where η_1 , η_2 be two constants in \mathbb{C} . Let us take $p_1(\vec{z_2}) = L(\vec{z_2}) + H(s) + W_1$, $p_2(\vec{z_2}) = L(\vec{z_2}) + H(s) + W_2$, where W_1 , W_2 are constants in \mathbb{C} , H(s) is a polynomial in $s = c_2 z_1 - c_1 z_2$.

Then combining this with (4.13) and (4.14), we obtain

$$b_{10} (d_1 + H'(s)c_2) + b_{01} (d_2 - H'(s)c_1) + b_{11} \{-H''(s)c_1c_2 - (d_1 + H'(s)c_2) (d_2 - H'(s)c_1)\} + b_{20} \{H''(s)c_2^2 - (d_1 + H'(s)c_2)^2\} + b_{02} \{H''(s)c_1^2 - (d_2 - H'(s)c_1)^2\} e^{L(\vec{c_2}) + W_1 - W_2} \equiv i, b_{10} (d_1 + H'(s)c_2) + b_{01} (d_2 - H'(s)c_1) + b_{11} \{-H''(s)c_1c_2 - (d_1 + H'(s)c_2) (d_2 - H'(s)c_1)\} + b_{20} \{H''(s)c_2^2 - (d_1 + H'(s)c_2)^2\} + b_{02} \{H''(s)c_1^2 - (d_2 - H'(s)c_1)^2\} e^{L(\vec{c_2}) - W_1 + W_2} \equiv i, b_{10} (d_1 + H'(s)c_2) + b_{01} (d_2 - H'(s)c_1) + b_{11} \{-H''(s)c_1c_2 + (d_1 + H'(s)c_2) (d_2 - H'(s)c_1)\} + b_{20} \{H''(s)c_1^2 + (d_2 - H'(s)c_1)^2\} e^{-L(\vec{c_2}) - W_1 + W_2} \equiv i, b_{10} \{H''(s)c_1^2 + (d_2 - H'(s)c_1)^2\} e^{-L(\vec{c_2}) - W_1 + W_2} \equiv i,$$

$$b_{10} (d_1 + H'(s)c_2) + b_{01} (d_2 - H'(s)c_1)$$

$$+ b_{11} \{-H''(s)c_1c_2 + (d_1 + H'(s)c_2) (d_2 - H'(s)c_1)\}$$

$$+ b_{20} \{H''(s)c_2^2 + (d_1 + H'(s)c_2)^2\}$$

$$+ b_{02} \{H''(s)c_1^2 + (d_2 - H'(s)c_1)^2\} e^{-L(\vec{c_2}) + W_1 - W_2} \equiv i.$$

Proceeding with the similar methods as done in Case 1 we conclude that $H(s) \equiv 0$. Then from (4.21) we obtain

$$-D_1(d_1, d_2)e^{L(\vec{c_2}) + W_1 - W_2} = i,$$

$$-D_1(d_1, d_2)e^{L(\vec{c_2}) - W_1 + W_2} = i,$$

$$D_2(d_1, d_2)e^{-L(\vec{c_2}) - W_1 + W_2} = i,$$

$$D_2(d_1, d_2)e^{-L(\vec{c_2}) + W_1 - W_2} = i.$$

Clearly we have

$$D_1(d_1, d_2)D_2(d_1, d_2) = 1, \quad e^{2L(\vec{c_2})} = -\frac{D_2(d_1, d_2)}{D_1(d_1, d_2)},$$
$$e^{2(W_1 - W_2)} = 1, \quad e^{W_1 - W_2} = -\frac{i}{D_1(d_1, d_2)}e^{-L(\vec{c_2})}.$$

In this case the form of the solution is

$$\begin{split} f_1(\vec{z_2}) &= \frac{e^{L(\vec{z_2}) - L(\vec{c_2}) + W_2} - e^{-L(\vec{z_2}) + L(\vec{c_2}) - W_2}}{2i}, \\ f_2(\vec{z_2}) &= \frac{e^{L(\vec{z_2}) - L(\vec{c_2}) + W_1} - e^{-L(\vec{z_2}) + L(\vec{c_2}) - W_1}}{2i}. \end{split}$$

5. Discussion related to Theorem 2.4 and an open question

From the expressions of $D_1(d_1,d_2)$, $D_2(d_1,d_2)$ we see that they are related in a certain way. More elaborately, when we consider only the second degree homogeneous differential operator, then $D_1(d_1,d_2)=D_2(d_1,d_2)$ and when we consider only the first order differential operator, then $D_1(d_1,d_2)=-D_2(d_1,d_2)$. Now we discuss the following cases:

Case 1: Let $D_1(d_1, d_2) = D_2(d_1, d_2) = D(d_1, d_2)$. Then from Case 1 and Case 4 in Theorem 2.4, we obtain $D^2(d_1, d_2) = 1$, that is $D(d_1, d_2) = \pm 1$.

Case 2: Let $D_1(d_1, d_2) = -D_2(d_1, d_2) = D(d_1, d_2)$. Then from Case 1 and Case 4 in Theorem 2.4, we obtain $D^2(d_1, d_2) = -1$, that is $D(d_1, d_2) = \pm i$.

Combining Case 1 and Case 2 we clearly see that $D_1(d_1, d_2)$ and $D_2(d_1, d_2)$ can take the values $\{1, -1, i, -i\}$ with $D_1(d_1, d_2)D_2(d_1, d_2) = 1$. In particular, we can write the solution of equation (2.2) as the follows: Let

$$f_1(\vec{z}) = \frac{S_{11}e^{-L(\vec{z_2}) - W_1} + S_{12}e^{L(\vec{z_2}) + W_1}}{2}, \quad f_2(\vec{z}) = \frac{S_{21}e^{-L(\vec{z_2}) - W_1} + S_{22}e^{L(\vec{z_2}) + W_1}}{2},$$

where W_1 , W_2 and S_{11} , S_{12} , S_{21} , S_{22} are constants in \mathbb{C} .

Now under the conclusion (A) in Theorem 2.4, we have the following:

(i)
$$D_1(d_1, d_2) = i$$
, $D_2(d_1, d_2) = -i$ and $e^{L(\vec{c_2})} = i$, $e^{W_1 + W_2} = -i$, then $S_{11} = -i$, $S_{12} = i$, $S_{21} = -1$, $S_{22} = -1$ or $e^{L(\vec{c_2})} = -i$, $e^{W_1 + W_2} = i$, then $S_{11} = -i$, $S_{12} = i$, $S_{21} = 1$, $S_{22} = 1$; or

- (ii) $D_1(d_1, d_2) = -i$, $D_2(d_1, d_2) = i$ and $e^{L(\vec{c_2})} = i$, $e^{W_1 + W_2} = i$, then $S_{11} = i$, $S_{12} = -i$, $S_{21} = -1$, $S_{22} = -1$ or $e^{L(\vec{c_2})} = -i$, $e^{W_1 + W_2} = -i$, then $S_{11} = i$, $S_{12} = -i$, $S_{21} = 1$, $S_{22} = 1$; or
- (iii) $D_1(d_1, d_2) = 1$, $D_2(d_1, d_2) = 1$ and $e^{L(\vec{c_2})} = 1$, $e^{W_1 + W_2} = i$, then $S_{11} = 1$, $S_{12} = 1$, then $S_{21} = i$, $S_{22} = -i$; $e^{L(\vec{c_2})} = -1$, $e^{W_1 + W_2} = -i$, then $S_{11} = 1$, $S_{12} = 1$, then $S_{21} = i$, $S_{22} = -i$; or
- (iv) $D_1(d_1, d_2) = -1$, $D_2(d_1, d_2) = -1$ and $e^{L(\vec{c_2})} = 1$, $e^{W_1 + W_2} = -i$, then $S_{11} = -1$, $S_{12} = -1$, $S_{21} = i$, $S_{22} = -i$ or $e^{L(\vec{c_2})} = -1$, $e^{W_1 + W_2} = i$, then $S_{11} = -1$, $S_{12} = -1$, $S_{21} = -i$, $S_{22} = i$.

Similarly, under conclusion (B) in Theorem 2.4, we have the following

- (i) $D_1(d_1, d_2) = -i$, $D_2(d_1, d_2) = i$, and $e^{L(\vec{c_2})} = 1$, $e^{W_1 W_2} = 1$, then $S_{11} = i$, $S_{12} = -i$, $S_{21} = i$, $S_{22} = -i$; or $e^{L(\vec{c_2})} = -1$, $e^{W_1 W_2} = -1$, then $S_{11} = i$, $S_{12} = -i$, $S_{21} = -i$, $S_{22} = i$; o
- (ii) $D_1(d_1, d_2) = i$, $D_2(d_1, d_2) = -i$ and $e^{L(\vec{c_2})} = 1$, $e^{W_1 W_2} = -1$, then $S_{11} = -i$, $S_{12} = i$, $S_{21} = i$, $S_{22} = -i$; or $e^{L(\vec{c_2})} = -1$, $e^{W_1 W_2} = 1$, then $S_{11} = -i$, $S_{12} = i$, $S_{21} = -i$, $S_{22} = i$; or
- (iii) $D_1(d_1, d_2) = -1$, $D_2(d_1, d_2) = -1$ and $e^{L(\vec{c_2})} = i$, $e^{W_1 W_2} = 1$, then $S_{11} = -1$, $S_{12} = -1$, $S_{21} = -1$, $S_{22} = -1$, or $e^{L(\vec{c_2})} = -i$, $e^{W_1 W_2} = -1$, then $S_{11} = -1$, $S_{12} = -1$, $S_{21} = 1$, $S_{22} = 1$; or (iv) $D_1(d_1, d_2) = 1$, $D_2(d_1, d_2) = 1$, and $e^{L(\vec{c_2})} = i$, $e^{W_1 W_2} = -1$, then $S_{11} = 1$,
- (iv) $D_1(d_1, d_2) = 1$, $D_2(d_1, d_2) = 1$, and $e^{L(\vec{c_2})} = i$, $e^{W_1 W_2} = -1$, then $S_{11} = 1$, $S_{12} = 1$, $S_{21} = -1$, $S_{22} = -1$; $e^{L(\vec{c_2})} = -i$, $e^{W_1 W_2} = 1$, then $S_{11} = 1$, $S_{12} = 1$, $S_{21} = 1$, $S_{22} = 1$.

In view of (2.1) and (2.2) the following question is inevitable:

What will be the possible form of transcendental entire solution of the following system of equation in \mathbb{C}^n

$$(P_{L_2}(f_1(\vec{z_2})))^2 + f_2(\vec{z_2} + \vec{c_2})^2 = Q_1(\vec{z_2}),$$

$$(P_{L_2}(f_2(\vec{z_2}))^2 + f_1(\vec{z_2} + \vec{c_2})^2 = Q_2(\vec{z_2});$$

where $Q_j(\vec{z_2}), j = 1, 2$ are two non-zero polynomials in \mathbb{C}^n ?

Acknowledgements. The authors would like to thank the anonymous referee for his/her valuable suggestions taht improved the overall presentation of the manuscript. J. Sarkar wishes to thank the Council of Scientific and Industrial Research (India), for their financial help under File No.-09/0106(13572)/2022-EMR-I.

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Abhijit Banerjee

Department of Mathematics, University of Kalyani, West Bengal 741235, India $Email\ address$: abanerjee_kal@yahoo.co.in

JHUMA SARKAR

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF KALYANI, WEST BENGAL 741235, INDIA Email address: jhumasarkar928@gmail.com