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EXISTENCE OF WEAK SOLUTIONS FOR NONLOCAL DIRICHLET PROBLEMS VIA YOUNG MEASURE THEORY

MOUAD ALLALOU, MOHAMED EL OUAARABI, ABDERRAHMANE RAJI

Abstract. This article investigates the existence of weak solutions for a class of nonlocal problems with Dirichlet boundary conditions. The proof of the existence result relies on Galerkin's approximation and Young's measure theory.

1. INTRODUCTION

Let D be a bounded domain in \mathbb{R}^n with smooth boundary $\partial \mathcal{D}$, and $1 < p < \infty$. We prove the existence of weak solutions for the nonlocal problem with Dirichlet boundary conditions,

$$
-g\left(\int_{\mathcal{D}} \mathcal{E}(v) \, \mathrm{d}z\right) \operatorname{div}[a(z, \nabla v) + |\nabla v|^{p-2} \nabla v] + |v|^{p-2} v = f(z, v) \quad \text{in } \mathcal{D},
$$

$$
v = 0 \quad \text{on } \partial \mathcal{D}, \tag{1.1}
$$

where

$$
\mathcal{E}(v) = \int_{\mathcal{D}} \left(\mathcal{A}(z, \nabla v) + \frac{1}{p} |\nabla v|^p \right) dz.
$$

The functions $f: \mathcal{D} \times \mathbb{R} \to \mathbb{R}, g, \mathcal{A}: \mathcal{D} \times \mathbb{R}^m \to \mathbb{R}, \text{ and } a: \mathcal{D} \times \mathbb{R}^m \to \mathbb{R}^m \text{ are }$ subject to conditions specified below.

The study of nonlinear boundary value problems has garnered significant attention over the past few decades, driven by advancements in fields such as elastic mechanics, electrorheological fluids, and image restoration; see [\[1,](#page-10-0) [7,](#page-10-1) [9,](#page-10-2) [17,](#page-11-0) [20,](#page-11-1) [23\]](#page-11-2).

Transmission problems appear in various applications in physics and biology (see [\[4,](#page-10-3) [8,](#page-10-4) [16\]](#page-11-3)). Recently, in [\[15\]](#page-11-4), the authors investigated the existence of ground-state solutions for a class of Kirchhoff-type transmission problems.

This work aims to explore the existence of weak solutions to problem [\(1.1\)](#page-0-0) by employing the principles of Young measures theory. Notably, our problem cannot be addressed with a variational framework because of the specific functions g and f. These functions introduce significant technical challenges, necessitating the use of alternative tools, such as Young measures, which facilitate the identification of weak limits. To the best of our knowledge, this is the first study to approach problem [\(1.1\)](#page-0-0) using this theoretical framework. For an exploration of closely related topics, readers are encouraged to consult references [\[10,](#page-10-5) [12\]](#page-10-6) and additional sources

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Young measures; Galerkin's approximation; Sobolev space.

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cited therein. For a thorough discussion on the steady-state case employing Young measures theory, we refer to [\[2,](#page-10-7) [3,](#page-10-8) [13,](#page-11-5) [22\]](#page-11-6).

A weak solution for [\(1.1\)](#page-0-0) is defined as a function $v \in W_0^{1,p}(\mathcal{D})$ satisfying the following equation for all $\Phi \in W_0^{1,p}(\mathcal{D}),$

$$
g\left(\int_{\mathcal{D}} \mathcal{F}(v) \,dz\right) \int_{\mathcal{D}} \left(a(z, \nabla v) + |\nabla v|^{p-2} \nabla v\right) \cdot \nabla \Phi \,dz + \int_{\mathcal{D}} |v|^{p-2} v \cdot \Phi \,dz
$$

=
$$
\int_{\mathcal{D}} f(z, v) \Phi \,dz.
$$

In this article we use the following assumptions:

(A1) $a: \mathcal{D} \times \mathbb{R}^m \to \mathbb{R}^m$ and $f: \mathcal{D} \times \mathbb{R} \to \mathbb{R}$ are Carathéodory functions which implies their measurability with respect to $z \in \mathcal{D}$ and continuity with respect to the other variables. Additionally, the mapping $\xi \mapsto a(z, \xi)$ is both a $C¹$ function and monotonic, i.e.,

$$
(a(z,\xi) - a(z,\xi')) \cdot (\xi - \xi') \ge 0 \quad \forall \xi, \xi' \in \mathbb{R}^m.
$$

(A2) We can find elements $\alpha_1, \alpha_2 \in L^{p'}(\mathcal{D})$, along with $\alpha_3 \in L^1(\mathcal{D})$ and positive constants $c_0, c_1 > 0$, such that

$$
|a(z,\xi)| \le c_0(\alpha_1(z) + |\xi|^{p-1}),
$$

\n
$$
|f(z,w)| \le \alpha_2(z) + |w|^q,
$$

\n
$$
p\mathcal{A}(z,\xi) \ge a(z,\xi) \cdot \xi \ge c_1 |\xi|^p - \alpha_3(z)
$$

for all $w \in \mathbb{R}$ and $0 \leq q < p - 1$.

- (A3) $A: \mathcal{D} \times \mathbb{R}^m \to \mathbb{R}$ is a Carathéodory function within the context of (A1). Furthermore, the mapping $\xi \mapsto \mathcal{A}(z, \xi)$ is both convex and C¹-function and it fulfills the relation $a(z,\xi) = \nabla \xi \mathcal{A}(z,\xi) = (\partial \mathcal{A}/\partial \xi)(z,\xi)$.
- (A4) $g: W^{1,p}(\mathcal{D}) \to (0, +\infty)$ are continuous and bounded on any bounded subset of $W^{1,p}(\mathcal{D})$ such that there are constants $g_0, g_1 > 0$ satisfying

$$
g_0 \le g(s) \le g_1.
$$

Our main result in this article reads as follows.

Theorem 1.1. Assume that $(A1)$ – $(A4)$ hold. Then, problem (1.1) has a weak solution in $W_0^{1,p}(\mathcal{D})$.

The structure of this paper is as follows: Section 2 offers a concise overview of essential aspects of Young measures. In Section 3, we focus on developing the approximate solutions and establishing preliminary estimates. The final section addresses various convergence outcomes and outlines the demonstration of the primary theorem.

2. Preliminaries

2.1. Fundamentals of Young measures. In this section, we provide a succinct summary of the fundamental concepts behind generalized Young measures and revisit relevant findings that will be employed in subsequent discussions. Our approach is influenced by the works of [\[18,](#page-11-7) [14\]](#page-11-8), with additional insights available in [\[19\]](#page-11-9) for a more thorough introduction.

We denote by $C_0(\mathbb{R}^m)$ the closure of the space comprising continuous functions on \mathbb{R}^m with compact support concerning the $|\cdot|_{\infty}$ -norm. Its dual space can be

identified as $\mathcal{M}(\mathbb{R}^m)$ the space encompassing signed Radon measures with finite mass. The duality pairing for $\sigma : \mathcal{D} \to \mathcal{M}(\mathbb{R}^m)$ is defined as

$$
\langle \sigma, \psi \rangle = \int_{\mathbb{R}^m} \psi(\eta) \, \mathrm{d}\sigma(\eta).
$$

Lemma 2.1 ([\[19\]](#page-11-9)). Assume that the sequence $\{y_\mu\}_{\mu \geq 1}$ is bounded in $L^\infty(\mathcal{D}; \mathbb{R}^m)$. Then there exist a subsequence still denoted $\{y_\mu\}_\mu$ and a Borel probability measure σ_z on \mathbb{R}^m for a.e. $z \in \mathcal{D}$, such that for almost each $h \in \mathcal{C}(\mathbb{R}^m)$ we have

$$
\psi(y_\mu) \to^* \bar{\psi} \quad weakly in \ L^\infty(\mathcal{D}),
$$

where

$$
\bar{\psi}(z) = \int_{\mathbb{R}^m} \psi(\eta) d\sigma_z(\eta).
$$

Definition 2.2. We call $\sigma = {\{\sigma_z\}_{z \in \mathcal{D}}}$ the family of Young measures associated with the subsequence $\{y_\mu\}_\mu$. It is shown in [\[6\]](#page-10-9), that if for all $R > 0$

$$
\limsup_{\mu \to \infty} |\{z \in \mathcal{D} \cap B_R(0) : |y_\mu(z)| \ge L\}| = 0,
$$

then for any measurable $\mathcal{D}' \subset \mathcal{D}$,

$$
\psi(z, y_{\mu}) \to \langle \sigma_z, \psi(z, .) \rangle = \int_{\mathbb{R}^m} \psi(z, \eta) d\sigma_z(\eta) \quad \text{weakly in } L^1(\mathcal{D})
$$

for any Carathéodory function $\psi : \mathcal{D}' \times \mathbb{R}^m \to \mathbb{R}$ such that $\psi(z, y_{\mu})$ is equiintegrable.

If we consider $y_{\mu} = \nabla w_{\mu}$, where $w_{\mu} : \mathcal{D} \to \mathbb{R}$, the above properties remain true, and the following lemma can be proved in a similar way as in [\[5,](#page-10-10) Lemma 4.1].

Lemma 2.3. Let (∇w_μ) be a bounded sequence in $L^p(\mathcal{D}; \mathbb{R}^m)$. Then the Young measure σ_z generated by ∇w_μ in $L^p(\mathcal{D};\mathbb{R}^m)$ satisfies:

- (1) $|\sigma_z|_{\mathcal{M}(\mathbb{R}^m)} = 1$ for a.e. $z \in \mathcal{D}$, i.e., σ_z is a probability measure.
- (2) The weak L^1 -limit of ∇w_μ is given by $\langle \sigma_z, id \rangle = \int_{\mathbb{R}^m} \eta \cdot d\sigma_z(\eta)$.
- (3) σ_z satisfies $\langle \sigma_z, id \rangle = \nabla v(z)$ for a.e. $z \in \mathcal{D}$.

We will need the following Fatou-type inequality.

Lemma 2.4. Let $\psi : \mathcal{D} \times \mathbb{R}^m \to \mathbb{R}$ be a Carathéodory function and $w_{\mu} : \mathcal{D} \to \mathbb{R}$ a sequence of measurable functions such that ∇w_{μ} generates the Young measure σ_z with $|\sigma_z|_{g(\mathbb{R}^m)} = 1$ for almost every $z \in \mathcal{D}$. Then

$$
\liminf_{\mu \to \infty} \int_{\mathcal{D}} \psi(z, \nabla w_{\mu}) dz \ge \int_{\mathcal{D}} \int_{\mathbb{R}^m} \psi(z, \eta) d\sigma_z(\eta) dz
$$

provided that the negative part of $\psi(z, \nabla w_\mu)$ is equi-integrable.

3. Proof of Theorem [1.1](#page-1-0)

Let us consider the functional $\mathcal{L}(v) : W_0^{1,p}(\mathcal{D}) \to \mathbb{R}$ given by

$$
\Phi \mapsto g \Big(\int_{\mathcal{D}} \left(\mathcal{A}(z, \nabla v) + \frac{1}{p} |\nabla v|^p \, \mathrm{d}z \right) \Big) \Big[\int_{\mathcal{D}} a(z, \nabla v) \cdot \nabla \Phi \, \mathrm{d}z \n+ \int_{\mathcal{D}} |\nabla v|^{p-2} \nabla v \cdot \nabla \Phi \, \mathrm{d}z \Big] + \int_{\mathcal{D}} |v|^{p-2} v \cdot \Phi \, \mathrm{d}z - \int_{\mathcal{D}} f(z, v) \Phi \, \mathrm{d}z,
$$

for arbitrary $v \in W_0^{1,p}(\mathcal{D})$ and $\Phi \in W_0^{1,p}(\mathcal{D})$.

Lemma 3.1. The functional $\mathcal{L}(v)$ is well defined, linear and bounded.

Proof. Firstly, utilizing Hölder inequality and conditions $(A2)$ – $(A4)$, we establish

$$
\begin{aligned} |\Lambda_1| &:= \Big|g\Big(\int_{\mathcal{D}} (A(z,\nabla v) + \frac{1}{p}|\nabla v|^p) \,\mathrm{d}z\Big) \\ &\quad \times \Big[\int_{\mathcal{D}} a(z,\nabla v)\nabla \Phi \,\mathrm{d}z + \int_{\mathcal{D}} |\nabla v|^{p-2}\nabla v\nabla \Phi \,\mathrm{d}z\Big]\Big| \\ &\le g_1\Big(\int_{\mathcal{D}} |a(z,\nabla v)|\cdot |\nabla \Phi| \,\mathrm{d}z + \int_{\mathcal{D}} |\nabla u|^{p-1}\cdot |\nabla \Phi| \,\mathrm{d}z\Big) \\ &\le g_1\Big(\int_{\mathcal{D}} c_0(\alpha_1(z) + |\nabla v|^{p-1})|\nabla \Phi| \,\mathrm{d}z\Big) + \|\nabla v\|_p^{p-1} \, \|\nabla \Phi\|_p \\ &\le C(\|\alpha_1\|_{p'} + |\nabla v\|_p^{p-1})\|\nabla \Phi\|_p + \|\nabla v\|_p^{p-1}\|\nabla \Phi\|_p \\ &\le C\|\nabla \Phi\|_p. \end{aligned}
$$

Conversely, we can also infer, based on the growth condition of f in $(A2)$ and the Hölder inequality, that

$$
\begin{aligned} |\Lambda_2| &:= \big| \int_{\mathcal{D}} f(z, v) \Phi \, dz \big| \\ &\le \int_{\mathcal{D}} |f(z, v) \Phi| \, dz \\ &\le (\|\alpha_2\|_{p'} + \|v\|_p^{p-1}) \|\Phi\|_p \\ &\le \lambda (\|\alpha_2\|_{p'} + \lambda^{p-1} \|\nabla v\|_p^{p-1}) \|\nabla \Phi\|_p, \end{aligned}
$$

with λ denoting the constant in Poincare's inequality, there exists a positive constant λ such that

$$
\|\Phi\|_{p} \le \lambda \|\nabla \Phi\|_{p} \quad \forall \Phi \in W_{0}^{1,p}(\mathcal{D}).
$$
\n(3.1)

On the other hand,

$$
|\Lambda_3| := |\int_{\mathcal{D}} |v|^{p-2} v \Phi \,dz|
$$

\n
$$
\leq \int_{\mathcal{D}} |v|^{p-1} |\Phi| \,dz
$$

\n
$$
\leq (\frac{1}{p'} + \frac{1}{p}) ||v||_p^{p-1} ||\nabla \Phi||_p.
$$

As the estimates of Λ_i for $i = 1, 2, 3$ are finite, $\mathcal{L}(v)$ is well defined. Moreover, $\mathcal{L}(v)$ is linear and for all $\Phi \in W_0^{1,p}(\mathcal{D})$, the inequality

$$
|\langle \mathcal{L}(v), \Phi \rangle| \le |\Lambda_1| + |\Lambda_2| + |\Lambda_3| \le C \|\nabla \Phi\|_p
$$

holds, indicating that $\mathcal{L}(v)$ is bounded.

By Lemma [3.1,](#page-3-0) we can define the operator $\mathcal{L}: W_0^{1,p}(\mathcal{D}) \to W^{-1,p'}(\mathcal{D})$, that satisfies the following result.

Proposition 3.2. The restriction of $\mathcal L$ to a finite dimensional linear subspace $\mathfrak D$ of $W_0^{1,p}(\mathcal{D})$ is continuous.

Proof. Let \mathfrak{O} be a finite linear subspace of $W_0^{1,p}(\mathcal{D})$. Suppose (v_μ) is a sequence in $\mathfrak D$ that converges to v in $\mathfrak D$.

Firstly, $v_{\mu} \to v$ and $\nabla v_{\mu} \to \nabla v$ almost everywhere.

$$
\sqcup
$$

$$
\int_{\mathcal{D}} |v_{\mu} - v|^p \, \mathrm{d}z \to 0 \quad \text{and} \quad \int_{\mathcal{D}} |\nabla v_{\mu} - \nabla v|^p \, \mathrm{d}z \to 0,
$$

since $v_{\mu} \to v$ strongly in \mathfrak{O} . Hence, there exist $Q_1, Q_2 \in L^1(\mathcal{D})$ such that $|v_{\mu}-v|^p \leq$ Q_1 and $|\nabla v_\mu - \nabla v|^p \leq Q_2$. We know that for $\gamma > 1$

$$
|t_1 + t_2|^{\gamma} \le 2^{\gamma - 1} (|t_1|^{\gamma} + |t_2|^{\gamma}).
$$

Then

$$
|v_{\mu}|^{p} = |v_{\mu} - v + v|^{p} \le 2^{p-1} (|v_{\mu} - v|^{p} + ||v||^{p}) \le 2^{p-1} (Q_{1} + ||v||^{p}).
$$

Like in the demonstration of $|v_\mu|^p$, it follows that $|v_\mu|^p$ and $|\nabla v_\mu|^p$ are bounded by a constant C. Thus, the continuity condition in (A1), (A3) and (A4) permits to deduce that

$$
g\Big(\int_{\Omega}(\mathcal{A}(z,\nabla v_k)+\frac{1}{p}|\nabla v_k|^p)dz\Big)\Big(\int_{\Omega}a(x,\nabla v_k)\nabla\Phi(z)\,\mathrm{d}z
$$

$$
-\int_{\Omega}|\nabla v_k|^{p-2}\nabla v_k\nabla\Phi(z)\,\mathrm{d}z\Big)+\int_{\Omega}|v_k|^{p-2}v_k\,\Phi\,\mathrm{d}z
$$

converges to

$$
g\Big(\int_{\Omega}(\mathcal{A}(z,\nabla v)+\frac{1}{p}|\nabla v|^p)dz\Big)\Big(\int_{\Omega}a(x,\nabla v)\nabla\Phi(z)\,\mathrm{d} z
$$

$$
-\int_{\Omega}|\nabla v|^{p-2}\nabla v\nabla\Phi(z)\,\mathrm{d} z\Big)+\int_{\Omega}|v|^{p-2}v.\Phi\,\mathrm{d} z
$$

and

$$
f(z, v_{\mu})\Phi(z) \to f(z, v)\Phi(z)
$$

almost everywhere as $k \to \infty$. Indeed, if \mathcal{D}' is a measurable subset of \mathcal{D} , and $\Phi \in W_0^{1,p}(\mathcal{D}), \text{ then}$

$$
\int_{\Omega'} |a(z, \nabla v_k) \cdot \nabla \Phi - |\nabla v_k|^{p-2} \nabla v_k \cdot \nabla \Phi| dz
$$
\n
$$
\leq \int_{\Omega'} c_0(\alpha_1(z) + |\nabla v_k|^{p-1}) |\nabla \Phi| dz + \int_{\Omega'} |\nabla v_k|^{p-1} \cdot |\nabla \Phi| dz
$$
\n
$$
\leq (c_0 |\alpha_1|_{p'} + (c_0 + 1) \underbrace{\|\nabla v_k\|_p^{p-1}}_{\leq C}) \Big(\int_{\Omega'} |\nabla \Phi|^p dz \Big)^{1/p}
$$

and (without loss of generality, we can assume $q = p - 1$)

$$
\int_{\mathcal{D}'} |f(z,v_{\mu})\Phi(z)| dz \leq \int_{\mathcal{D}'} (\alpha_2(z) + |v_{\mu}|^{p-1}) |\Phi| dz
$$

$$
\leq \lambda (|\alpha_2|_{p'} + \underbrace{||v_{\mu}||_p^{p-1}}_{\leq C}) \Big(\int_{\mathcal{D}'} |\nabla \Phi|^p dz \Big)^{1/p},
$$

Using Hölder's and Poincaré inequalities, along with (3.1) . Moreover, we have

$$
g(\int_{\mathcal{D}'}\left(\mathcal{A}(z,\nabla v_{\mu})+\frac{1}{p}|\nabla v_{\mu}|^{p}\right)\mbox{d} z)\leq g_{1}<\infty,
$$

by (A4) and the boundedness of $|\nabla v_\mu|_p$. By utilizing the Vitali Theorem, we can establish the continuity of \mathcal{L} . \Box **Remark 3.3.** In this section, we have used only the condition $q \leq p-1$. Thus Lemma [3.1](#page-3-0) and Proposition [3.2](#page-3-2) are still valid as $q = p - 1$.

Now, the problem [\(1.1\)](#page-0-0) is equivalent to find a solution $v \in W_0^{1,p}(\mathcal{D})$ such that

$$
\langle \mathcal{L}(v), \Phi \rangle = 0 \quad \text{for all } \Phi \in W_0^{1,p}(\mathcal{D}).
$$

To find such a solution we apply a Galerkin scheme. Since $W_0^{1,p}(\mathcal{D})$ is separable there exists a sequence (\mathfrak{O}_{μ}) of finite dimensional subspaces such that $\cup_{\mu\geq 1}\mathfrak{O}_{\mu}$ is dense in $W_0^{1,p}(\mathcal{D})$. Let $\{x_1,\ldots,x_r\}$ be a basis of \mathfrak{O}_μ where $\dim \mathfrak{O}_\mu = r$. Next, Let us define

$$
\mathcal{G}: \mathbb{R}^r \to \mathbb{R}^r,
$$

$$
(d_i)_{i=1,\ldots,r} \to (\langle \mathcal{L}(d_i x_i), x_j \rangle)_{j=1,\ldots,r}.
$$

Proposition 3.4. $\mathcal L$ is continuous and $\mathcal G(d) \cdot d \to \infty$ as $|d|_{\mathbb R^r} \to \infty$.

Proof. G is trivially continuous, by the continuity of $\mathcal L$ restricted to \mathfrak{O}_{μ} (see Propo-sition [3.2](#page-3-2) if necessary). Consider $d \in \mathbb{R}^r$ and $v = d_i x_i \in \mathfrak{O}_{\mu}$ (with conventional summation). The condition $|d|_{\mathbb{R}^r} \to \infty$ is equivalent to $||v||_{1,p} \to \infty$, and we have

$$
\mathcal{G}(d) \cdot d = \langle \mathcal{L}(v), v \rangle.
$$

Note that

$$
\Lambda_4 := g \Big(\int_{\mathcal{D}} (A(z, \nabla v) + \frac{1}{p} |\nabla v|^p) dz \Big) \Big[\int_{\mathcal{D}} a(x, \nabla v) \cdot \nabla v dz + \int_{\mathcal{D}} |\nabla v|^p dz \Big]
$$

\n
$$
\ge g_0 \Big[\int_{\mathcal{D}} a(x, \nabla v) \cdot \nabla v dz + \int_{\mathcal{D}} |\nabla v|^p dz \Big] \quad \text{(by (A4))}
$$

\n
$$
\ge g_0 \Big(\int_{\mathcal{D}} c_1 |\nabla v|^p dz - \int_{\mathcal{D}} \alpha_3(z) dz \Big) + g_0 \int_{\mathcal{D}} |\nabla v|^p dz
$$

\n
$$
\ge C_{\min} \int_{\mathcal{D}} |\nabla v|^p dz - C' \int_{\mathcal{D}} \alpha_3(z) dz,
$$

since $\beta \geq 1$. Finally, from the growth condition (A2) and [\(3.1\)](#page-3-1) we have

$$
\begin{aligned} |\Lambda_5| &:= |\int_{\mathcal{D}} f(z, v)v \,dz| \\ &\leq \int_{\mathcal{D}} |f(z, v)v| \,dz \\ &\leq \int_{\mathcal{D}} (\alpha_2(z) + \|v\|^{p-1}) \|v\| \,dz \\ &\leq \lambda \|\alpha_2\|_{p'} \|\nabla v\|_p + \lambda^{p+1} \|\nabla v\|_p^{p+1}. \end{aligned}
$$

Hence

$$
\langle \mathcal{L}(v), v \rangle \ge \Lambda_4 - \Lambda_5
$$

\n
$$
\ge C_{min} \|\nabla v\|_p^p - C' \|\alpha_3\|_{p'} - \lambda \|\alpha_2\|_{p'} \|\nabla u\|_p - \lambda^{p+1} \|\nabla v\|_p^{p+1} \to \infty
$$

as $||v||_{1,p} \to \infty$, since C_{min} , $C' > 0$ and $(p > max(1, q + 1))$.

Proposition 3.5. For all $k \in \mathbb{N}$ there exists $v_{\mu} \in \mathfrak{O}_{\mu}$ such that

$$
\langle \mathcal{L}(v_{\mu}), \Phi \rangle = 0 \quad \text{for all } \Phi \in \mathfrak{O}_{\mu}.
$$

Proof. By Proposition [3.4,](#page-5-0) there exists $R > 0$ such that for all $d \in \partial B_R(0) \subset \mathbb{R}^r$ we have $\mathcal{G}(d) \cdot d > 0$, and the usual topological argument [\[24,](#page-11-10) Proposition 2.8], there exists $z \in B_R(0)$ such that $\mathcal{G}(z) = 0$. Hence, for all $k \in \mathbb{N}$ there exists $v_\mu \in \mathfrak{O}_\mu$ such that $\langle \mathcal{L}(v_\mu), \varphi \rangle = 0$ for all $\varphi \in \mathfrak{O}_\mu$.

Proposition 3.6. The constructed sequence (v_{μ}) in Proposition [3.5](#page-5-1) is uniformly bounded, i.e., there is a constant $R > 0$ such that $||v_\mu||_{1,p} \leq R$ for all $k \in \mathbb{N}$.

Proof. By Proposition [3.4](#page-5-0) there exists $R > 0$ with the property that $\langle \mathcal{L}(v), v \rangle > 1$ whenever $||v||_{1,p} > R$. Hence, for the sequence of Galerkin approximations $v_{\mu} \in \mathfrak{O}_{\mu}$ which satisfy $\langle \mathcal{L}(v_\mu), v_\mu \rangle = 0$ by the Proposition [3.5,](#page-5-1) we get the uniform boundedness of (v_μ) in $W_0^{1,p}$ $(\mathcal{D}).$

4. Proofs and properties for the convergence

In this section, we present general convergence findings pertaining to the functions denoted as $a(\cdot)$, $\mathcal{A}(\cdot)$, and $f(\cdot)$. Given that the sequence (v_μ) remains within bounded limits in the space $W_0^{1,p}(\mathcal{D})$, as established in Propositions [3.4,](#page-5-0) [3.5,](#page-5-1) and [3.6,](#page-6-0) we can infer, based on the assertions of Lemma [2.4,](#page-2-0) the existence of a Young measure denoted as σ_x . This measure is generated by the gradients of v_μ within the space $L^p(\mathcal{D};\mathbb{R}^m)$.

We define

$$
\widetilde{a}(z,\nabla v) = a(z,\nabla v) + |\nabla v|^{p-2} \nabla v,
$$

where \tilde{a} adheres to conditions (A1)-(A3) with both coercivity and growth rate set to p , specifically,

$$
\widetilde{a}(z,\xi).\xi \ge |\xi|^p,
$$

$$
|\widetilde{a}(z,\xi)| \le |\xi|^{p-1} + S(c_3,p,q).
$$
 (4.1)

Lemma 4.1. The Young measure σ_z generated by ∇v_μ satisfies

$$
(\widetilde{a}(z,\eta) - \widetilde{a}(z,\nabla v)) \cdot (\eta - \nabla v) = 0 \quad on \, \operatorname{supp} \sigma_z,
$$

where $\text{supp }\sigma_z$ is the support of σ_z for a.e. $z \in \mathcal{D}$.

Proof. We consider the sequence

$$
e_{\mu} := (\widetilde{a}(z, \nabla v_{\mu}) - \widetilde{a}(z, \nabla v)) \cdot (\nabla v_{\mu} - \nabla v)
$$

= $\widetilde{a}(z, \nabla v_{\mu}) \cdot (\nabla v_{\mu} - \nabla v) - \widetilde{a}(z, \nabla v) \cdot (\nabla v_{\mu} - \nabla v)$
= $e_{\mu,1} + e_{\mu,2}$.

Given the growth condition of \tilde{a} in [\(4.1\)](#page-6-1) and the weak convergence described in Lemma [2.3,](#page-2-1) it follows that

$$
\liminf_{\mu \to \infty} \int_{\mathcal{D}} e_{\mu,2} dz = \int_{\mathcal{D}} \widetilde{a}(z, \nabla \nu) \cdot \left(\underbrace{\int_{\mathbb{R}^m} \eta d\sigma_z(\eta)}_{=: \nabla \nu(z)} - \nabla \nu \right) dz = 0.
$$

Thus

$$
e := \liminf_{\mu \to \infty} \int_{\mathcal{D}} e_{\mu} dz = \liminf_{\mu \to \infty} \int_{\mathcal{D}} e_{\mu,1} dz.
$$

By the growth condition on \tilde{a} in [\(4.1\)](#page-6-1), $(\tilde{a}(z, \nabla v_{\mu}) \cdot \nabla v)$ is equi-integrable. Let us fix an arbitrary measurable subset $\mathcal{D}' \subset \mathcal{D}$. Then, the coercivity condition in [\(4.1\)](#page-6-1) implies

$$
\int_{\mathcal{D}'} |\min(\tilde{a}(z, \nabla v_{\mu}) \cdot \nabla v_{\mu}, 0)| \, dz \le \int_{\mathcal{D}'} |\nabla v_{\mu}|^p \, dz < \infty,
$$
\n(4.2)

This shows the equi-integrability of $(\widetilde{a}(z,\nabla v_{\mu})\cdot\nabla v_{\mu})$. Therefore $(\widetilde{a}(z,\nabla v_{\mu})\cdot(\nabla v_{\mu} \nabla v$)) is also equi-integrable, and by Lemma [2.4,](#page-2-0) this yields

$$
\int_{\mathcal{D}} \int_{\mathbb{R}^m} \widetilde{a}(z, \eta) \cdot (\eta - \nabla v) d\sigma_z(\eta) dz \le \liminf_{\mu \to \infty} \int_{\mathcal{D}} \widetilde{a}(z, \nabla v_{\mu}) \cdot (\nabla v_{\mu} - \nabla v) dz = e.
$$

Now, let us show that $e \leq 0$. By Propositions [3.5,](#page-5-1) we can write

$$
g\left(\int_{\Omega} (\mathcal{A}(z, \nabla v_k) + \frac{1}{p} |\nabla v_k|^p) dz\right)
$$

\$\times \left(\int_{\Omega} a(x, \nabla v_k) \cdot (\nabla v_k - \nabla v_k) dz - \int_{\Omega} |\nabla v_k|^{p-2} \nabla v_k \cdot (\nabla v_k - \nabla v) dz\right)\$
=
$$
\int_{\Omega} f(z, v_k)(v_k - v) dz - \int_{\Omega} |v_k|^{p-2} v_k.(v_k - v) dz.
$$

By (A4) we have

$$
g_0 \int_{\mathcal{D}} a(z, \nabla v_\mu) \cdot (\nabla v_\mu - \nabla v) dz
$$

\n
$$
\leq \int_{\mathcal{D}} f(z, v_\mu) (v_\mu - v) dz - \int_{\mathcal{D}} |v_\mu|^{p-2} v_\mu (v_\mu - v) dz
$$

\n
$$
+ g_0 \int_{\mathcal{D}} |\nabla v_\mu|^{p-2} \nabla v_\mu \cdot (\nabla v_\mu - \nabla v) dz.
$$

Then

$$
\int_{\mathcal{D}} a(z, \nabla v_{\mu}) \cdot (\nabla v_{\mu} - \nabla v) dz
$$
\n
$$
\leq \frac{1}{g_0} \Big(\int_{\mathcal{D}} f(z, v_{\mu})(v_{\mu} - v) dz - \int_{\mathcal{D}} |v_{\mu}|^{p-2} v_{\mu} (v_{\mu} - v) dz \Big)
$$
\n
$$
- \int_{\mathcal{D}} |\nabla v_{\mu}|^{p-2} \nabla v_{\mu} \cdot (\nabla v_{\mu} - \nabla v) dz
$$
\n
$$
\leq C_m \int_{\mathcal{D}} f(z, v_{\mu})(v_{\mu} - v) dz
$$
\n
$$
\leq C_m (\|d_2\|_{p'} + \underbrace{\|v_{\mu}\|_{p}^{p-1}}_{\leq C}) \|v_{\mu} - v\|_{p} \to 0 \quad \text{as } k \to \infty.
$$

This is achieved by Hölder's inequality and the fact that $v_{\mu} \to v$ in $W_0^{1,p}(\mathcal{D})$. Hence, we have

$$
\int_{\mathcal{D}} \int_{\mathbb{R}^m} \widetilde{a}(z,\eta) \cdot (\eta - \nabla v) d\sigma_z(\eta) dz \le 0.
$$

In conclusion, we can infer from this and Equation [\(4.2\)](#page-7-0) that

$$
\int_{\mathcal{D}} \int_{\mathbb{R}^m} \left(\widetilde{a}(z, \eta) - \widetilde{a}(z, \nabla v) \right) \cdot (\eta - \nabla v) \, \mathrm{d}\sigma_z(\eta) \, \mathrm{d}z \le 0.
$$

The function \tilde{a} being monotonic, the integral above evaluates to zero with respect to the product measure $d\sigma_z(\eta) \otimes dz$, meaning that

$$
\int_{\mathcal{D}} \int_{\mathbb{R}^m} \left(\widetilde{a}(z, \eta) - \widetilde{a}(z, \nabla v) \right) \cdot (\eta - \nabla v) d\sigma_z(\eta) \otimes dz = 0.
$$

Consequently, we obtain

$$
(\widetilde{a}(z,\eta) - \widetilde{a}(z,\nabla v)) \cdot (\eta - \nabla v) = 0 \text{ on } \text{supp } \sigma_z.
$$

Proposition 4.2. For almost every $z \in \mathcal{D}$, the support of σ_z is contained within the set where \tilde{A} coincides with the supporting hyper-plane L defined as

$$
L := \{ (\eta, \mathcal{A}(z, \nabla v) + \widetilde{a}(z, \nabla v) \cdot (\eta - \nabla v)) \},
$$

that is

$$
\operatorname{supp} \sigma_z \subset K_z = \{ \eta \in \mathbb{R}^m : \widetilde{\mathcal{A}}(z, \eta) = \widetilde{\mathcal{A}}(z, \nabla \nu) + \widetilde{a}(z, \nabla \nu) \cdot (\eta - \nabla \nu) \}.
$$

Proof. Let $\eta \in \text{supp } \sigma_z$. By Lemma [4.1](#page-6-2) implies for all $t \in [0,1]$, we have

$$
(1-t)\big(\widetilde{a}(z,\eta) - \widetilde{a}(z,\nabla v)\big) \cdot (\eta - \nabla v) = 0. \tag{4.3}
$$

Therefore, by the monotonicity condition and [\(4.3\)](#page-8-0), we obtain

$$
0 \le (1-t)\Big(\widetilde{a}(z,\eta) - \widetilde{a}(z,\nabla v + t(\eta - \nabla v))\Big) \cdot (\eta - \nabla v)
$$

= $(1-t)\Big(\widetilde{a}(z,\nabla v) - \widetilde{a}(z,\nabla v + t(\eta - \nabla v))\Big) \cdot (\eta - \nabla v).$ (4.4)

Using the monotonicity condition, we have

$$
(\widetilde{a}(z,\nabla v) - \widetilde{a}(z,\nabla v + t(\eta - \nabla v))) \cdot t(\nabla v - \eta) \ge 0,
$$

and since $t \in [0, 1]$, we deduce that

$$
(\tilde{a}(z,\nabla v) - \tilde{a}(z,\nabla v + t(\eta - \nabla v))) \cdot (1 - t)(\nabla v - \eta) \ge 0.
$$
 (4.5)

Combining [\(4.4\)](#page-8-1) and [\(4.5\)](#page-8-2) we find that

$$
(\widetilde{a}(z,\nabla v) - \widetilde{a}(z,\nabla v + t(\eta - \nabla v))\big)\cdot (\eta - \nabla v) = 0.
$$

It follows from (A3) that

$$
\widetilde{\mathcal{A}}(z,\eta) = \widetilde{\mathcal{A}}(z,\nabla\upsilon) + \int_0^1 \widetilde{a}(z,\nabla\upsilon + t(\eta - \nabla\upsilon)) \cdot (\eta - \nabla\upsilon) dz
$$

$$
= \widetilde{\mathcal{A}}(z,\nabla\upsilon) + \widetilde{a}(z,\nabla\upsilon) \cdot (\eta - \nabla\upsilon).
$$

Hence $\eta \in K_z$, i.e., supp $\sigma_z \subset K_z$ for almost every $z \in \mathcal{D}$.

Now, we establish the proof of our main result.

Proof of Theorem [1.1.](#page-1-0) Since $\xi \mapsto \widetilde{A}(z,\xi)$ is convex, we can represent it as

$$
\widetilde{\mathcal{A}}(z,\eta) =: \mathcal{H}(\eta) \ge \widetilde{\mathcal{A}}(z,\nabla \nu) + \widetilde{a}(z,\nabla \nu) \cdot (\eta - \nabla \nu) =: R(\eta)
$$

for all $\eta \in \mathbb{R}^m$. Assuming that $\eta \mapsto \mathcal{H}(\eta)$ is a C¹-function, as specified in the hypothesis, we obtain the following relationships for any $\xi \in \mathbb{R}^m, t \in \mathbb{R}$

$$
\frac{\mathcal{H}(\eta+t\xi)-\mathcal{H}(\eta)}{t} \ge \frac{R(\eta+t\xi)-R(\eta)}{t} \quad \text{for } t>0,
$$

$$
\frac{\mathcal{H}(\eta+t\xi)-\mathcal{H}(\eta)}{t}\leq \frac{R(\eta+\tau\xi)-R(\eta)}{t} \quad \text{for } t<0.
$$

Consequently, we can deduce that $\nabla_{\eta} \mathcal{H} = \nabla_{\eta} R$, which implies

$$
\widetilde{a}(z,\eta) = \widetilde{a}(z,\nabla\upsilon) \quad \text{for all } \eta \in K_z \supset \text{supp } \sigma_z. \tag{4.6}
$$

Since $\tilde{a}(z, \nabla v_{\mu})$ is equi-integrable, by [\(4.6\)](#page-9-0) and Lemma [2.3,](#page-2-1) its weak L^1 -limit satisfies isfies

$$
\bar{a}(z) = \int_{\mathbb{R}^m} \tilde{a}(z, \eta) d\sigma_z(\eta)
$$

\n
$$
= \int_{\text{supp }\sigma_z} \tilde{a}(z, \eta) d\sigma_z(\eta)
$$

\n
$$
= \int_{\text{supp }\sigma_z} \tilde{a}(z, \nabla v) d\sigma_z(\eta)
$$

\n
$$
= \tilde{a}(z, \nabla v).
$$
\n(4.7)

Next, if we consider the following Carathéodory function

$$
\mathcal{B}(z,\eta) = |\tilde{a}(z,\eta) - \bar{a}(z)|, \quad \eta \in \mathbb{R}^m,
$$

then, since $a(z, \nabla v_\mu)$ is equi-integrable, we conclude that $\mathcal{B}_\mu(z) := \mathcal{B}(z, \nabla v_\mu)$ is also equi-integrable, and its weak L^1 -limit is

$$
\mathcal{B}_{\mu}\to\bar{\mathcal{B}}\qquad\text{in }L^{1}(\mathcal{D}),
$$

where

$$
\bar{\mathcal{B}}(z) = \int_{\mathbb{R}^m} |\tilde{a}(z,\eta) - \bar{a}(z)| d\sigma_z(\eta) = \int_{\text{supp }\sigma_z} |\tilde{a}(z,\eta) - \bar{a}(z)| d\sigma_z(\eta) = 0,
$$

by [\(4.5\)](#page-8-2) and [\(4.6\)](#page-9-0). Notably, the convergence of \mathcal{B}_{μ} is strong since $\mathcal{B}_{\mu} \geq 0$. Applying [\(4.1\)](#page-6-1), we derive theinequalities

$$
\widetilde{\mathcal{A}}(z,\nabla v_{\mu}) \geq \frac{1}{p} \widetilde{a}(z,\nabla v_{\mu}) \cdot \nabla v_{\mu} \geq \frac{1}{p} |\nabla v_{\mu}|^{p};
$$

thus

$$
\int_{\mathcal{D}'} |\min(\widetilde{\mathcal{A}}(z,\nabla v_{\mu}),0)| \,dz < \infty.
$$

Consequently, $\mathcal{A}(z, \nabla v_\mu)$ is both bounded and equi-integrable. As a result, its weak L^1 -limit is $\int_{\mathbb{R}^m} \widetilde{\mathcal{A}}(z,\eta) d\sigma_z(\eta)$. Utilizing Lemma [4.1,](#page-6-2) we obtain

$$
\int_{\mathbb{R}^m} \widetilde{\mathcal{A}}(z,\eta) d\sigma_z(\eta) = \int_{\text{supp }\sigma_z} \widetilde{\mathcal{A}}(z,\eta) d\sigma_z(\eta) = \widetilde{\mathcal{A}}(z,\nabla \nu),
$$

by Equation [\(4.2\)](#page-7-0). The continuity of the function g in (A4) and $\mathcal{B}_{\mu} \to 0$ in $L^1(\mathcal{D})$, imply that

$$
g\left(\int_{\Omega} (\mathcal{A}(z, \nabla v_k) + \frac{1}{p} |\nabla v_k|^p) dz\right)
$$

\$\times \left(\int_{\Omega} a(z, \nabla v_k) \nabla \Phi(z) dz - \int_{\Omega} |\nabla v_k|^{p-2} \nabla v_k \nabla \Phi(z) dz\right) + \int_{\Omega} |v_k|^{p-2} v_k dz\$

converges to

$$
g\Big(\int_{\Omega} (A(z,\nabla v) + \frac{1}{p}|\nabla v|^p) \,\mathrm{d}z\Big)
$$

$$
\times \Big(\int_{\Omega} a(z, \nabla v) \nabla \Phi(z) dz - \int_{\Omega} |\nabla v|^{p-2} \nabla v \nabla \Phi(z) dz \Big) + \int_{\Omega} ||v||^{p-2} v dz.
$$

To complete the proof, we need to consider the term $\int_{\mathcal{D}} f(z, v_{\mu}) \Phi(z) dz$. We know that (v_{μ}) is bounded in $W_0^{1,p}(\mathcal{D})$ according to Propositions [3.4,](#page-5-0) [3.5,](#page-5-1) and [3.6,](#page-6-0) up to a subsequence, $v_{\mu} \to v$ in $L^p(\mathcal{D})$. For some ϵ positive, we have

$$
\int_{\mathcal{D}} |v_{\mu} - v|^p dz \ge \int_{\{z \in \mathcal{D}: |v_{\mu} - v| \ge \epsilon\}} |v_{\mu} - v|^p dz \ge \epsilon^p |\{z \in \mathcal{D}: |v_{\mu} - v| \ge \epsilon\}|,
$$

which implies

$$
|\{z \in \mathcal{D} : |v_{\mu} - v| \ge \epsilon\}| \le \frac{1}{\epsilon^p} \int_{\mathcal{D}} |v_{\mu} - v|^p dz \to 0 \text{ as } k \to \infty,
$$

thus $v_{\mu} \rightarrow v$ in measure and almost everywhere. The continuity of the function f in $(A1)$ implies

$$
f(z,\upsilon_\mu)\Phi(z)\to f(z,\upsilon)\Phi(z)
$$

almost everywhere. From the growth condition in $\mathcal F$ and the uniform bound in Propositions [3.4,](#page-5-0) [3.5,](#page-5-1) and [3.6,](#page-6-0) it follows that $f(z, v_\mu)\Phi(z)$ is equi-integrable. Consequently,

$$
f(z, v_\mu)\Phi(z) \to f(z, v)\Phi(z)
$$
 in $L^1(\mathcal{D})$,

and then, by the Vitali convergence theorem, we establish that

$$
\lim_{\mu \to \infty} \int_{\mathcal{D}} f(z, v_{\mu}) \Phi(z) dz = \int_{\mathcal{D}} f(z, v) \Phi(z) dz \quad \forall \Phi \in \bigcup_{\mu \ge 1} \mathfrak{O}_{\mu}.
$$

Since $\cup_{\mu\geq 1} \mathfrak{O}_{\mu}$ is dense in $W_0^{1,p}(\mathcal{D})$, it follows that v is a weak solution of [\(1.1\)](#page-0-0) as desired. \Box

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Mouad Allalou

Applied Mathematics and Scientific Computing Laboratory, Faculty of Science and Technics, Sultan Moulay Slimane University, Beni Mellal, BP 523, 23000, Morocco Email address: mouadallalou@gmail.com

Mohamed El Ouaarabi

FUNDAMENTAL AND APPLIED MATHEMATICS LABORATORY, FACULTY OF SCIENCES AÏN CHOCK, Hassan II University, Casablanca, BP 5366, 20100, Morocco.

Applied Mathematics and Scientific Computing Laboratory, Faculty of Science and Technics, Sultan Moulay Slimane University, Beni Mellal, BP 523, 23000, Morocco Email address: mohamed.elouaarabi@etu.univh2c.ma, mohamedelouaarabi93@gmail.com

Abderrahmane Raji

Applied Mathematics and Scientific Computing Laboratory, Faculty of Science and Technics, Sultan Moulay Slimane University, Beni Mellal, BP 523, 23000, Morocco Email address: rajiabd2@gmail.com